

# Magnetorotational Supernovae and Jet Formation

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**Supernova is one of the most powerful explosion in the Universe, energy (radiation and kinetic) about  $10^{51}$  egr**

**End of the evolution of massive stars, with initial mass more than about 8 Solar mass.**

Tracks in HR diagram of a representative selection of stars from the main sequence till the end of the evolution Iben (1985)

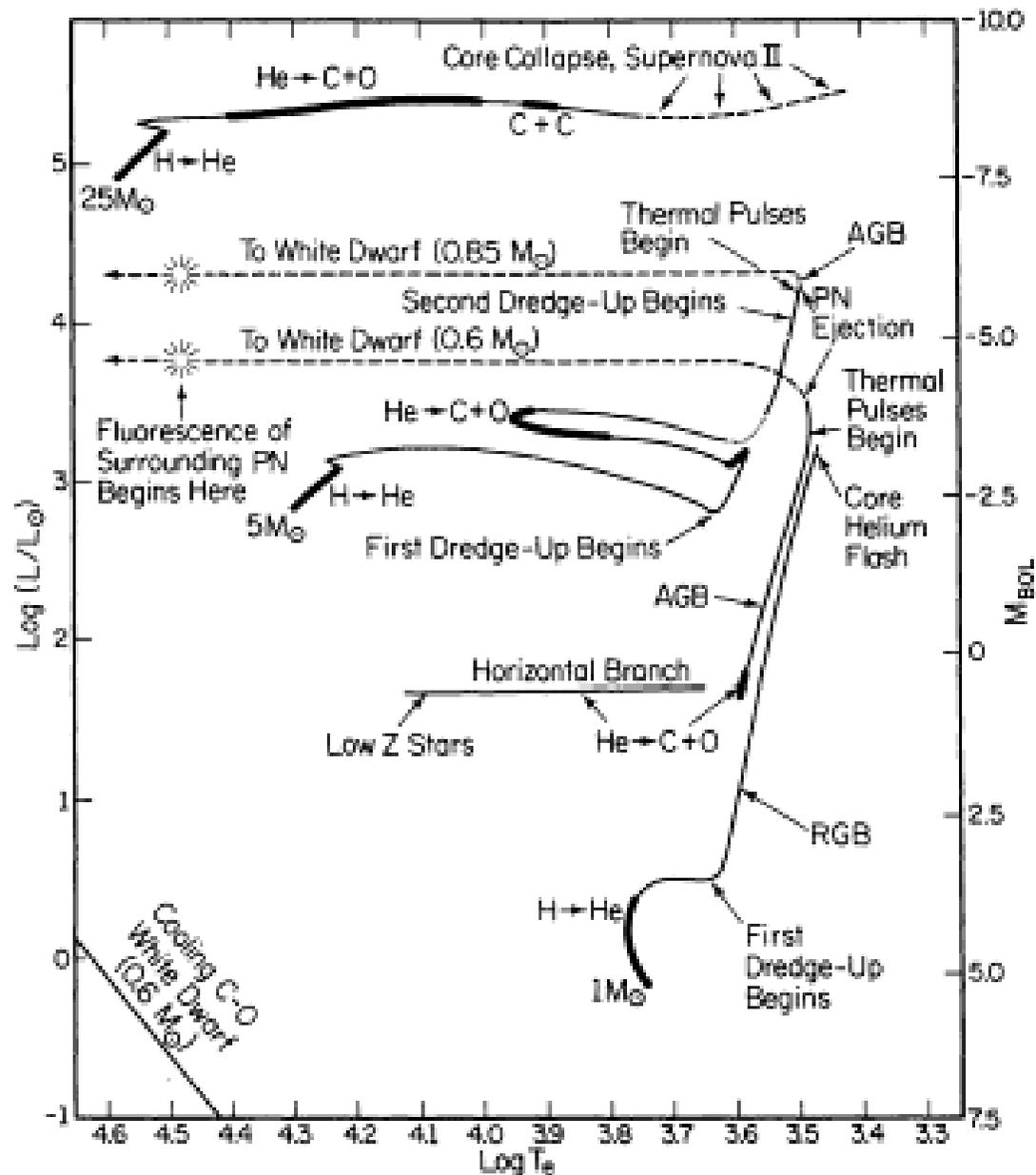


FIG. 15. Tracks in the HR diagram of a representative selection of stars. The heavy portions of each curve define locations where the major core nuclear burning phases occur. Details of tracks during transitory phases between major nuclear burning phases are suppressed. For stars of initial mass less than about  $2\text{-}3 M_{\odot}$ , where in the HR diagram subsequent core helium burning ( $\text{He} \rightarrow \text{C} + \text{O}$ ) takes place

# Explosion mechanisms of spherically symmetric star

1. **Thermonuclear explosion** of C-O degenerate core (**SN Ia**)

1. **Core collapse and formation of a neutron star, neutrino  
deposition**

gravitational energy release up to  $5 \cdot 10^{53}$  erg, carried away by  
neutrino (**SN II, SN Ib,c**)

**Equal to binding energy of the neutron star**

# NUCLEOSYNTHESIS IN SUPERNOVAE

F. HOYLE

St. John's College, Cambridge, and California Institute of Technology

AND

WILLIAM A. FOWLER

California Institute of Technology

*Received May 21, 1960*

## ABSTRACT

The role of Type I and Type II supernovae in nucleosynthesis is treated in some detail. It is concluded that  $s$ -process formation of the iron-group elements takes place in Type II supernovae, while  $r$ -process formation of the neutron-rich isotopes of the heavy elements takes place in Type I supernovae. The explosion of Type II supernovae is shown to follow implosion of the non-degenerate core material. The explosion of Type I supernovae results from the ignition of degenerate nuclear fuel in stellar material.

Astrophysical Journal, vol. 143, p.626 (1966)

## The Hydrodynamic Behavior of Supernovae Explosions

**S.Colgate, R.White**

*Received June 29, 1965*

We regard the release of gravitational energy attending a dynamic change in configuration to be the primary energy source in supernovae explosions. Although we were initially inspired by and agree in detail with the mechanism for initiating gravitational instability proposed by Burbidge, Burbidge, Fowler, and Hoyle, we find that the dynamical implosion is so violent that an energy many times greater than the available thermonuclear energy is released from the star's core and transferred to the star's mantle in a supernova explosion. The energy released corresponds to the change in gravitational potential of the unstable imploding core; the transfer of energy takes place by the emission and deposition of neutrinos.

D. Arnett (1966, 1967)

L.Ivanova, V. Imshennik, D. Nadyozhin (1969)

In a simple 1-D model neutrino deposition cannot give enough energy to matter (heating) for SN explosion

**Neutrino convection – Epstein (1979)**

leads to emission of higher energy neutrino, may transfer more energy into heating

**Results are still controversial**

2D and 3D calculations

3D calculations give less effective explosion because of opposite energy fluxes in 2 and 3 D  
**formation of larger eddies in 2D (artificial)**

**formation of smaller eddies in 3D (real),**

1968: PULSARS – rapidly rotating, strongly magnetized neutron stars

**Magnetorotational explosion (MRE):** transformation of the rotational energy of the neutron star into explosion energy by means of the magnetic field in core collapse SN

Soviet Astronomy, Vol. 14, p.652 (1971)

The Explosion of a Rotating Star As a Supernova Mechanism.

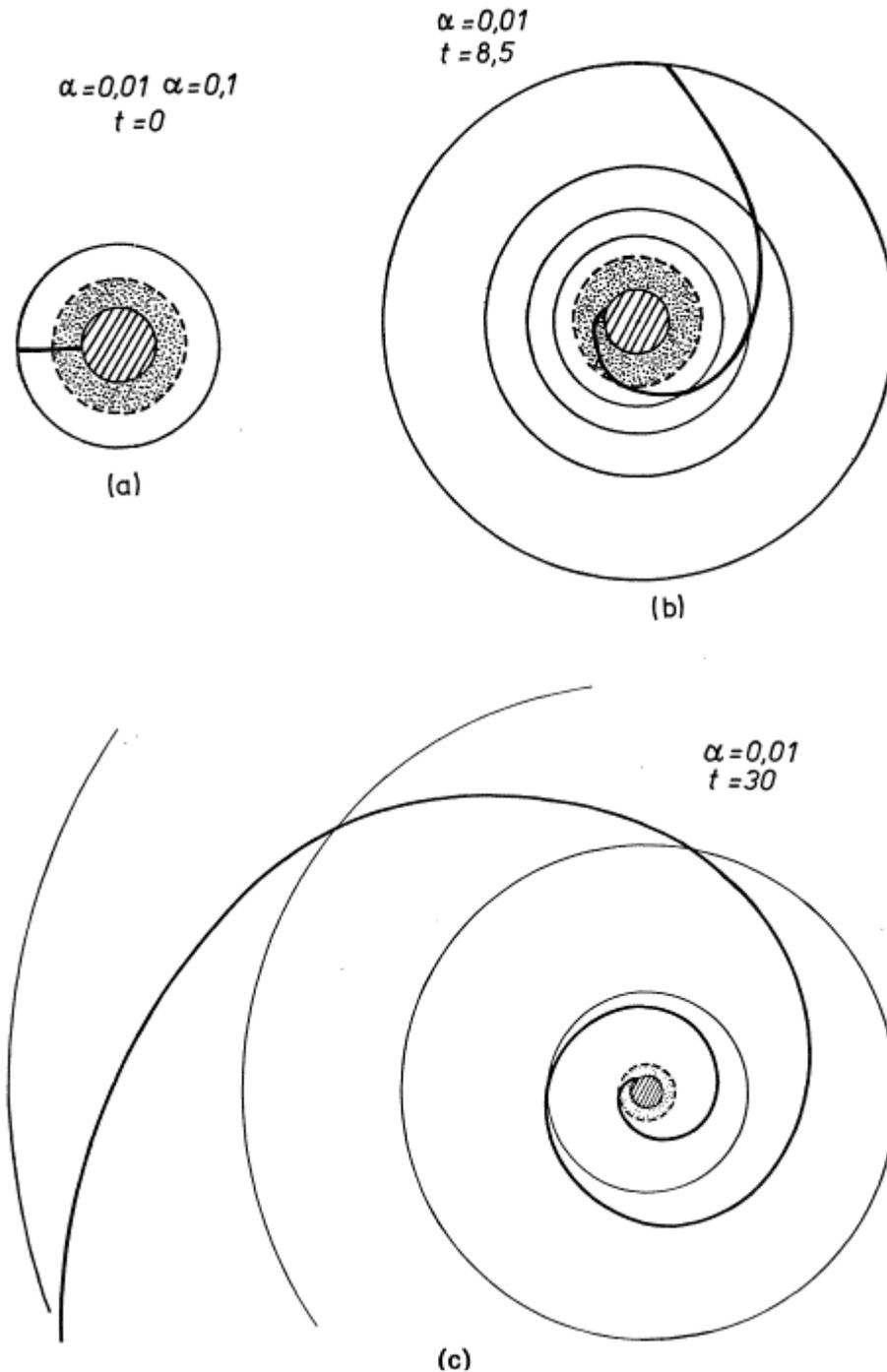
**G.S.Bisnovatyi-Kogan**

Translated from *Astronomicheskii Zhurnal*, Vol. 47, No. 4,  
pp. 813-816, July-August, 1970

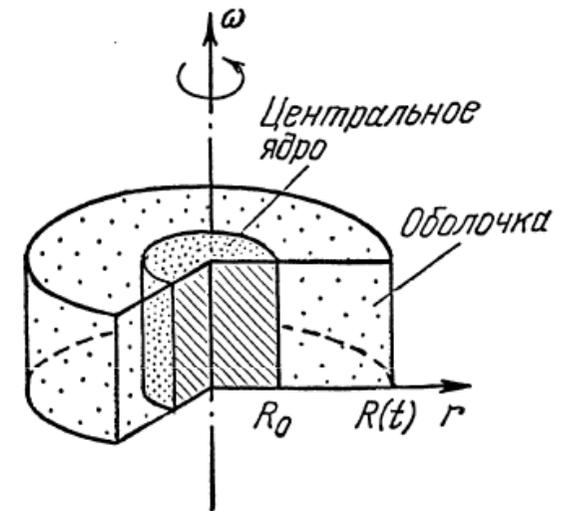
Original article submitted September 3, 1969

A supernova-explosion mechanism not associated with nuclear detonation is proposed for a rotating star. It involves the transfer of angular momentum from a rapidly rotating neutron star, formed through collapse, to the envelope, where the centrifugal force is nearly equal to the gravitational force. An explosion will result when the centrifugal force inside the envelope exceeds the gravitational force, with a shock wave being generated. Angular momentum will be transferred efficiently if a sufficiently strong magnetic field,  $H \approx 3 \cdot 10^9$  gauss, is present.

(1976)



$\alpha=0.01,$   
 $t=8.5$



1-D calculations of  
magnetorotational  
explosion

$$\alpha = \frac{E_{mag0}}{E_{grav}}$$

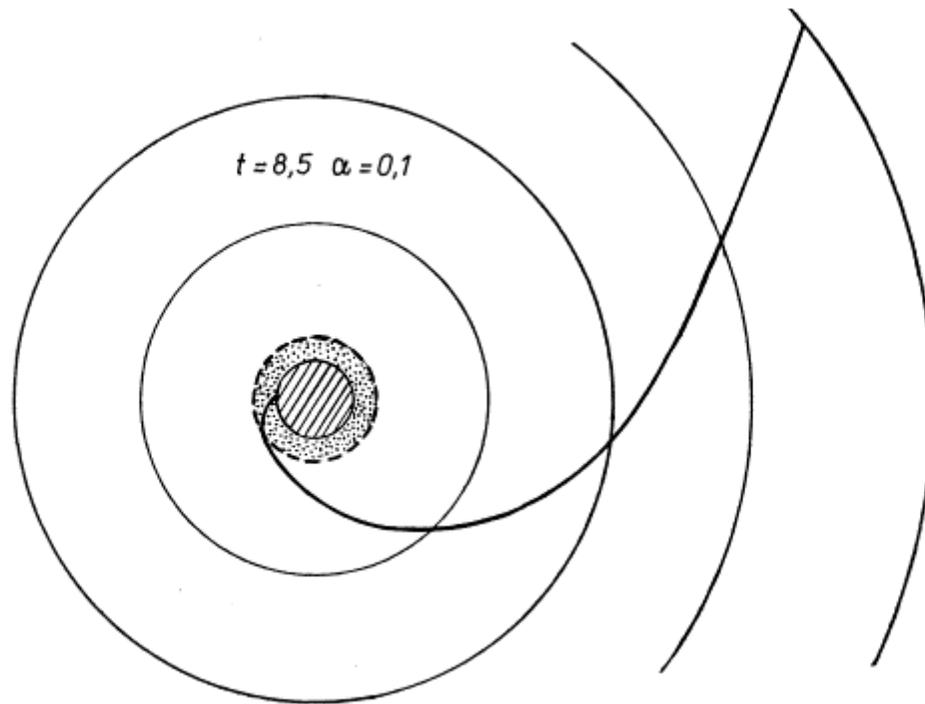
1-D calculations of  
magnetorotational explosion  
B.-K., Popov, Samokhin (1976).

Ardeljan, Bisnovatyi-Kogan, Popov (1979),  
Astron. Zh., **56**, 1244

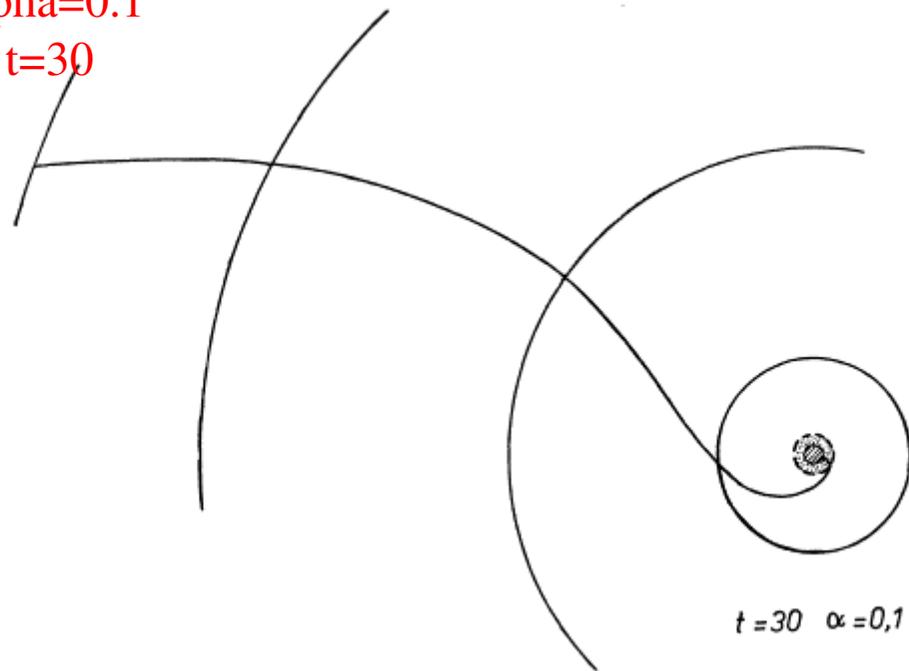
$\alpha=10^{-2}, 10^{-4}, 10^{-8}$

$\alpha=10^{-2}$ - dashed line,  $\alpha=10^{-4}$ - full line

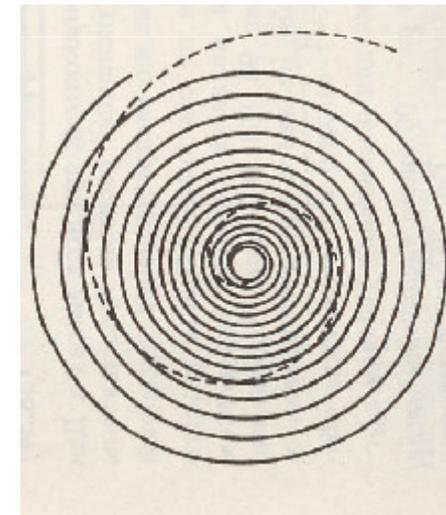
$\alpha=0.1$   
 $t=30$



(a)



(b)



## The main results of 1-D calculations:

**Magneto-rotational explosion (MRE) has an efficiency about 10% of rotational energy.**

For the neutron star mass the ejected mass  $\approx 0.1M_{\odot}$ ,

Explosion energy  $\approx 10^{51}$  erg

Ejected mass and explosion energy depend very weakly on the parameter  $\alpha$

**Explosion time strongly depends on  $\alpha$ .**

$$\text{Explosion time} = t_{\text{взрыва}} \sim \frac{1}{\sqrt{\alpha}}$$

Small  $\alpha$  is difficult for numerical calculations with EXPLICIT numerical schemes because of the Courant restriction on the time step, “stiff” system of equations:  $\alpha$  determines a “stiffness”.

**In 2-D numerical IMPLICIT schemes have been used.**

**Difference scheme** (Ardeljan, Chernigovskii, Kosmachevskii, Moiseenko)

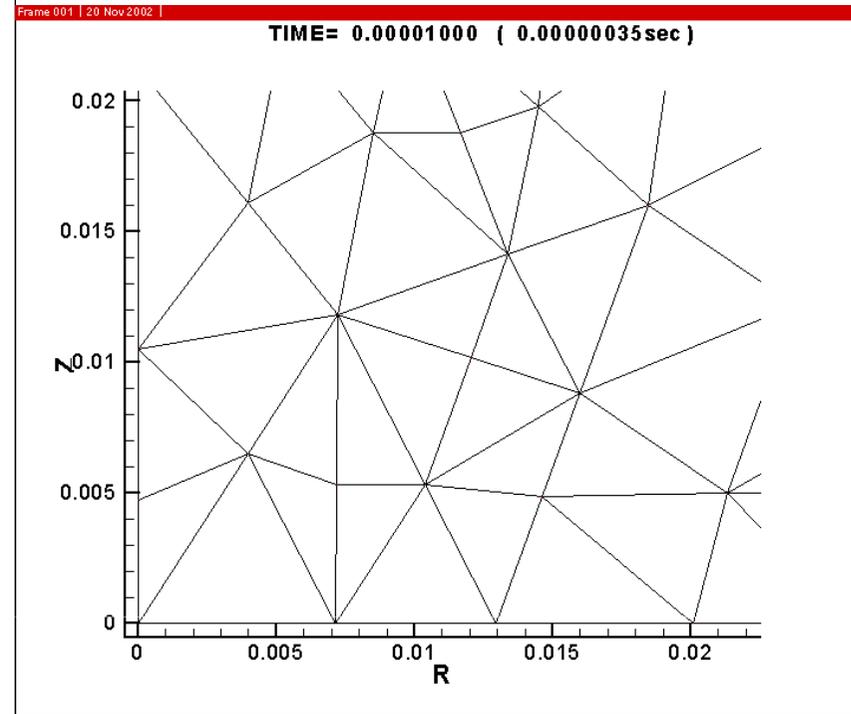
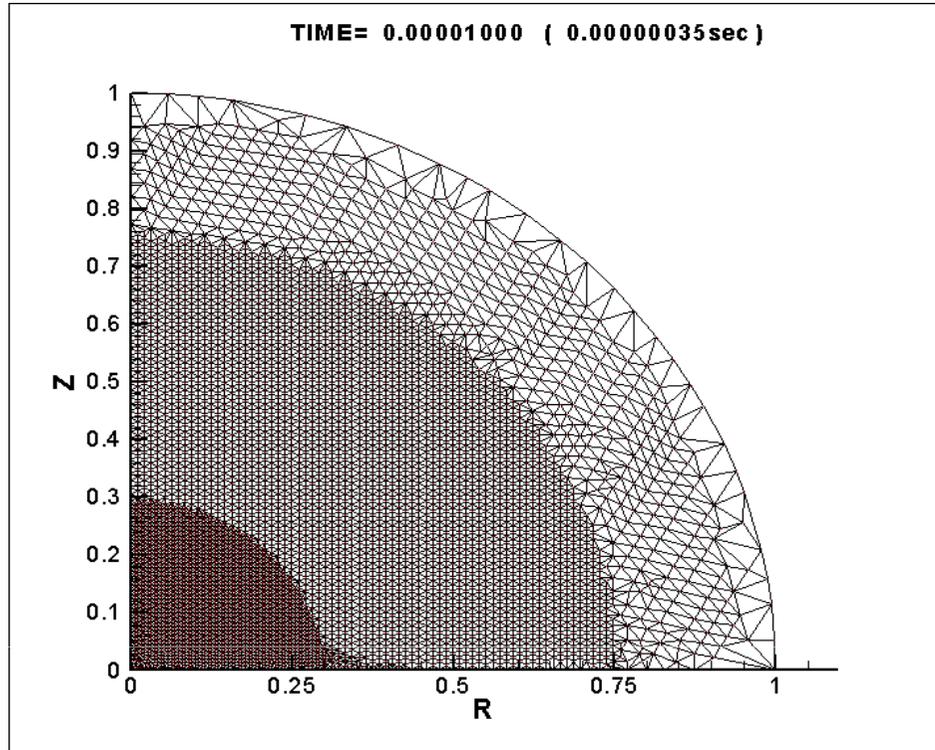
**Lagrangian, on triangular reconstructing grid, implicate, fully conservative**

**Ardeljan N.V., Kosmachevskii K.V., Chernigovskii S.V., 1987,** Problems of construction and research of conservative difference schemes for magneto-gas-dynamics, MSU, Moscow (in Russian)

**Ardeljan N.V., Kosmachevskii K.V. 1995,** Computational mathematics and modeling, **6**, 209

**Ardeljan N.V., Bisnovatyi-Kogan G.S., Kosmachevskii K.V., Moiseenko S.G., 1996,** Astron. Astrophys. Supl.Ser., **115**, 573

# Grid reconstruction (example)



# Presupernova Core Collapse

Ardeljan et. al., 2004, *Astrophysics*, **47**, 47

**Equations of state** takes into account degeneracy of electrons and neutrons, relativity for the electrons, nuclear transitions and nuclear interactions. Temperature effects were taken into account approximately by the addition of the pressure of radiation and of an ideal gas.

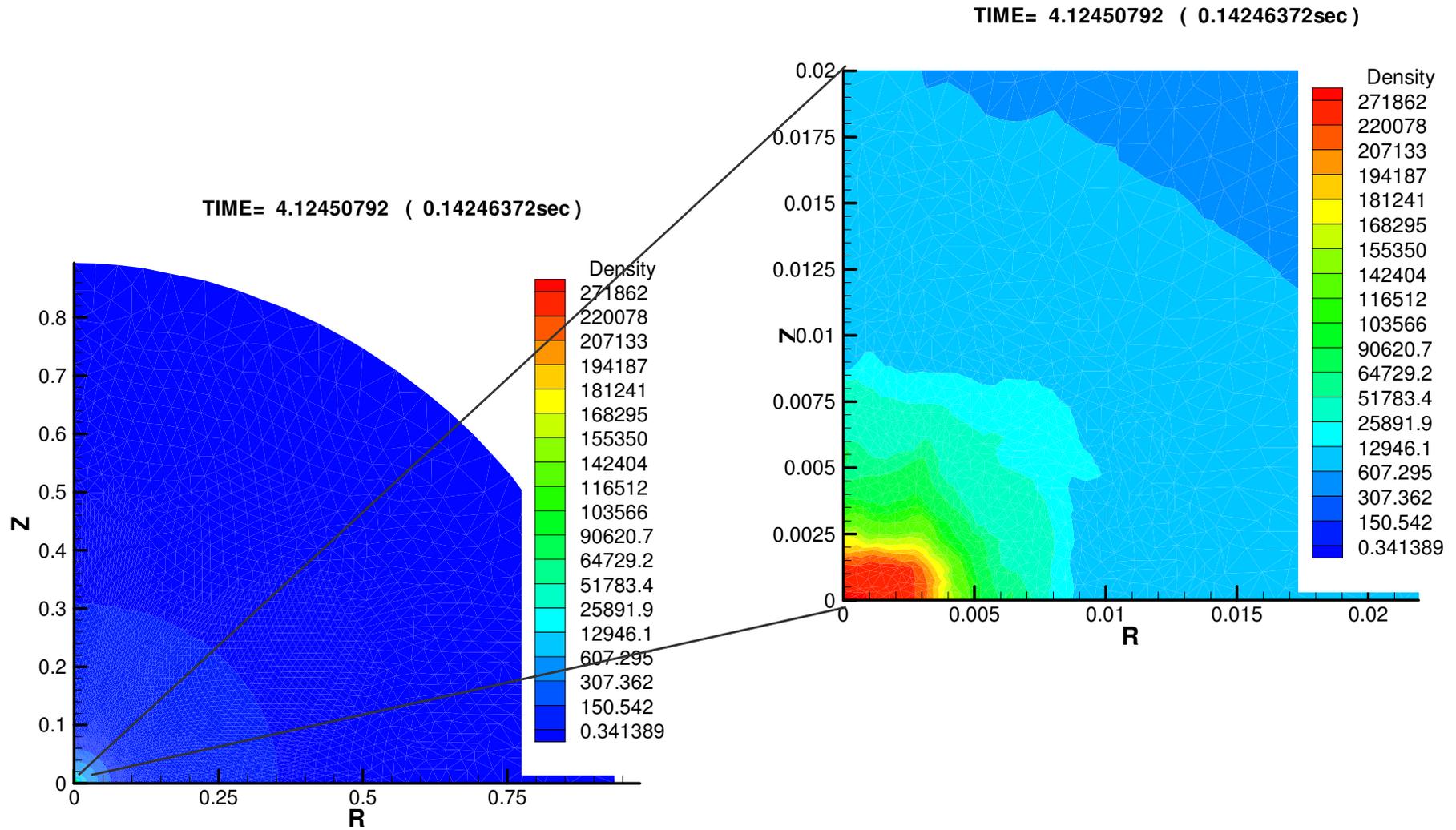
**Neutrino losses** and **iron dissociation** were taken into account in the energy equations.

A cool iron white dwarf was considered at the stability border with a mass equal to the Chandrasekhar limit.

$$M = 1.0042 \cdot M_{sun}$$

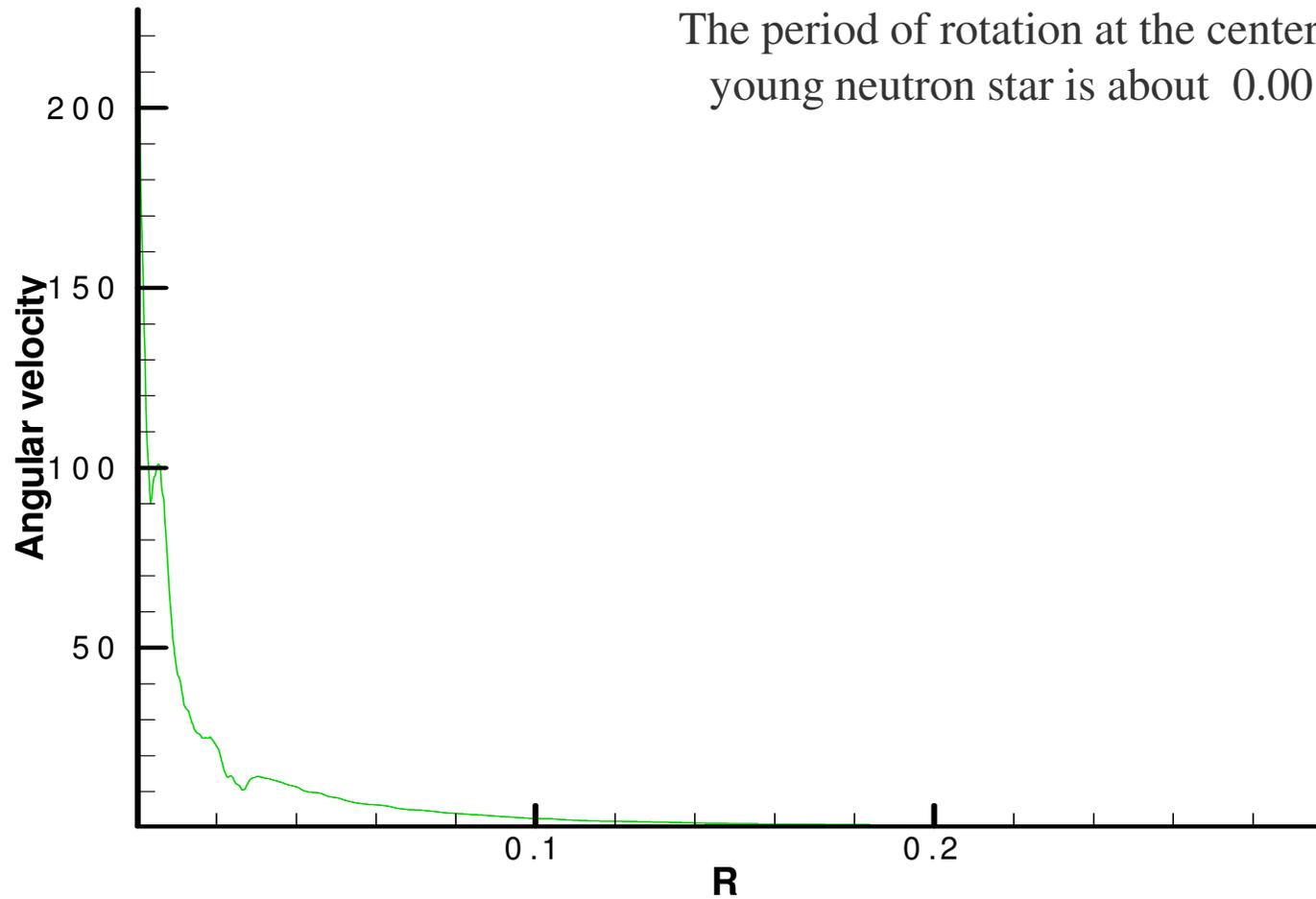
To obtain the collapse we increase the density at each point by 20% and switch on a uniform rotation.

# Maximal compression state



Shock wave does not produce SN explosion :

## Distribution of the angular velocity



## **2-D magnetorotational supernova**

N.V.Ardeljan, G.S.BK, S.G.Moiseenko MNRAS, **359**, 333 (2005)

## **A magnetorotational core-collapse model with jets**

S. G. Moiseenko, G. S. BK and N. V. Ardeljan MNRAS **370**, 501 (2006)

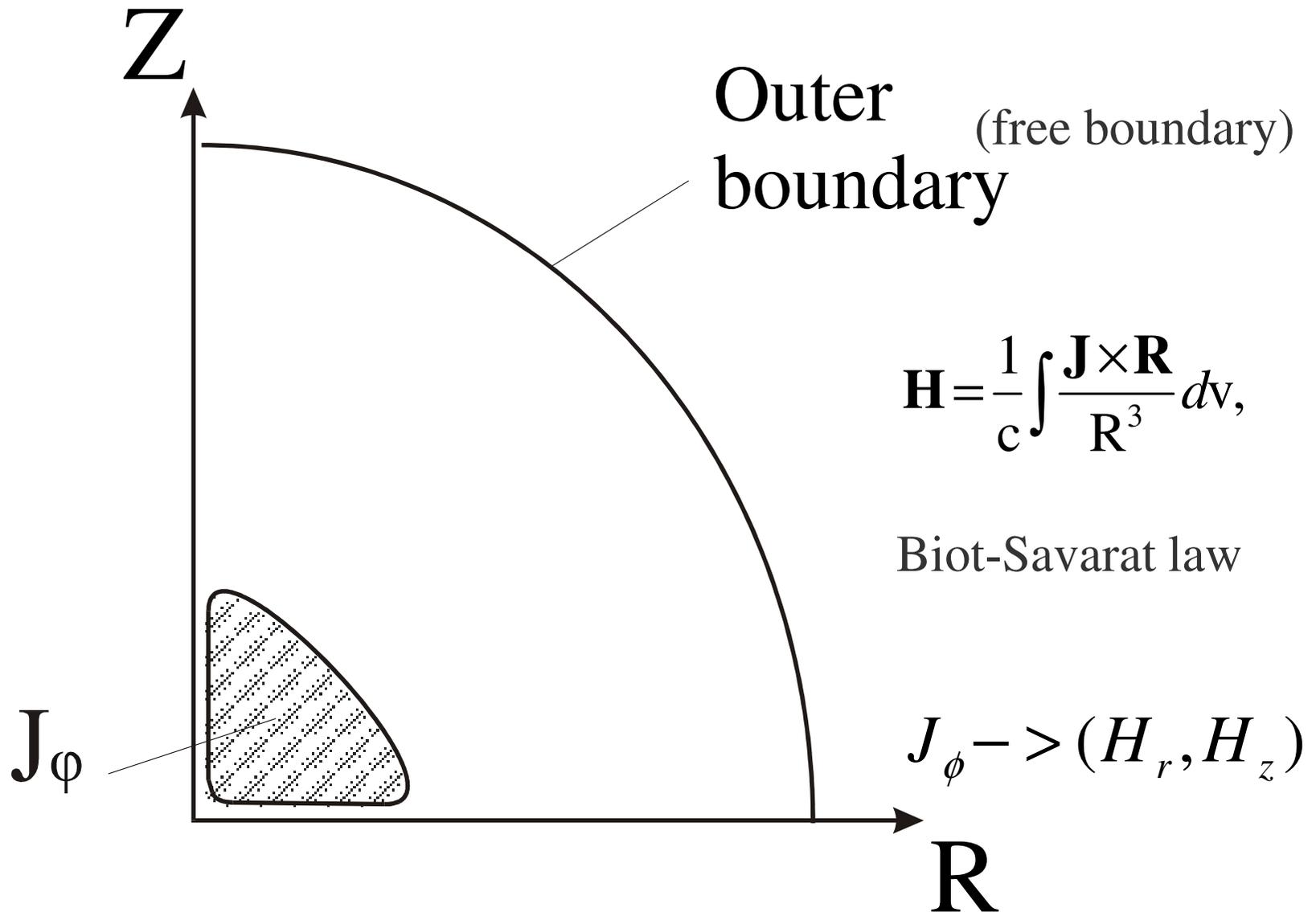
## **Different magneto-rotational supernovae**

BK, G. S.; Moiseenko, S. G.; Ardelyan, N. V., Astr. Rep.**52**, 997 (2008)

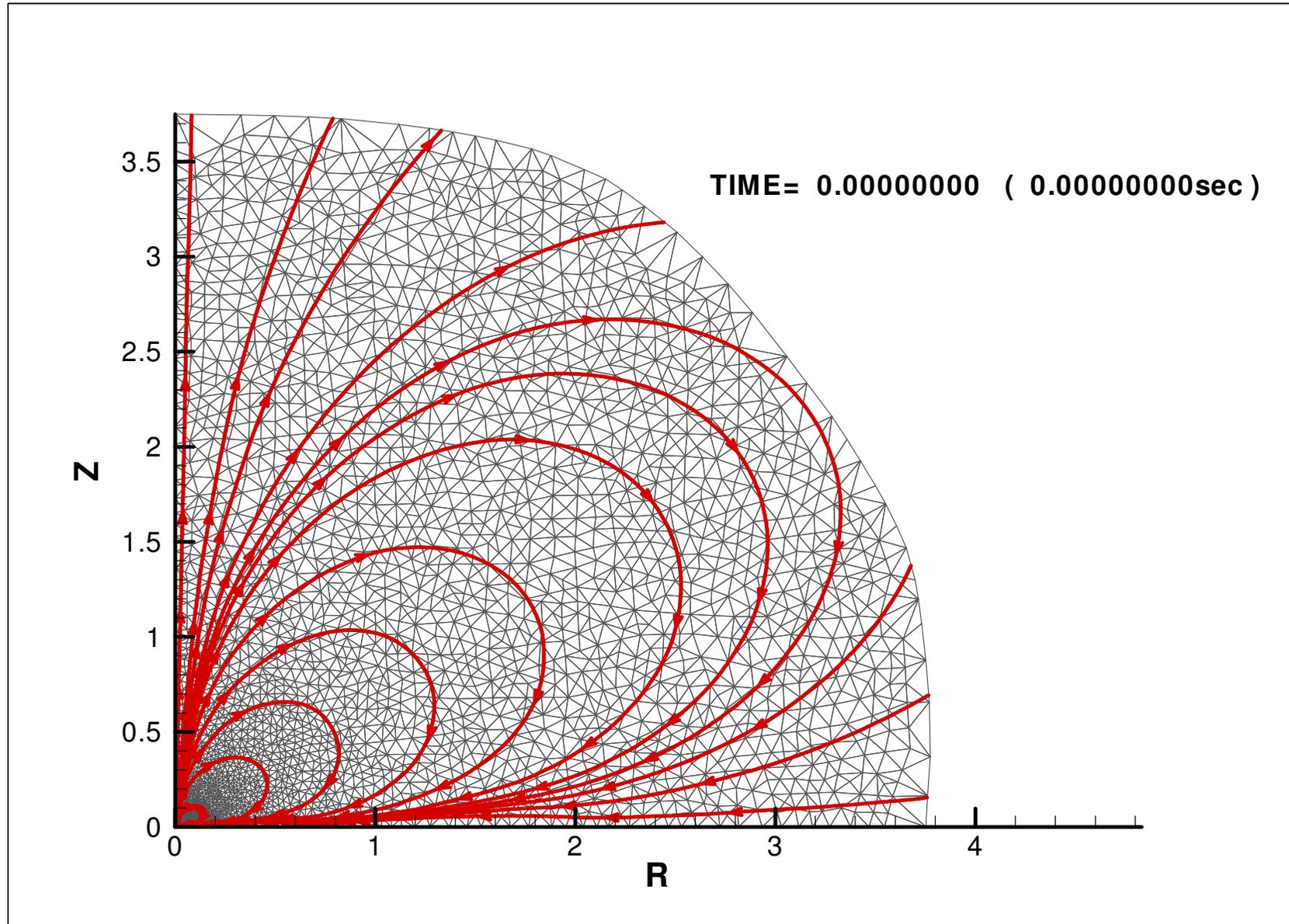
Equations: MHD + self-gravitation, infinite conductivity.

Axial symmetry (  $\frac{\partial}{\partial\phi} = 0$  ), equatorial symmetry (z=0).

Initial toroidal current  $J_\phi$



# Initial magnetic field –quadrupole-like symmetry

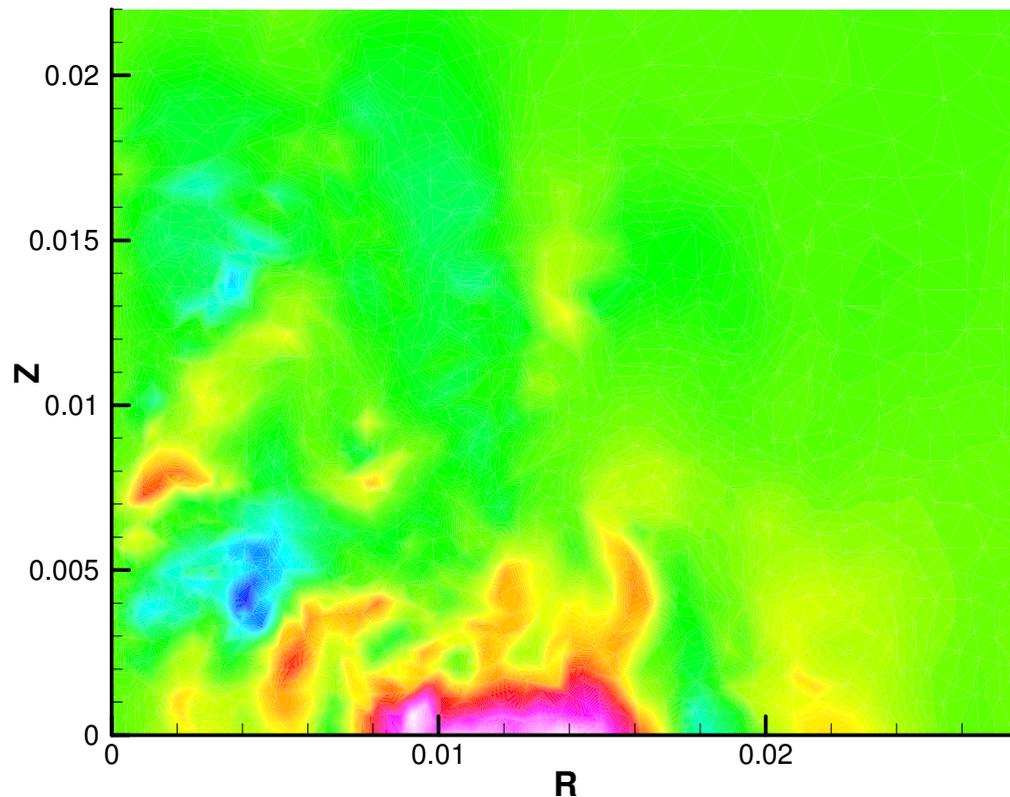


## Toroidal magnetic field amplification.

**pink** – maximum\_1 of  $H_f^2$     **blue** – maximum\_2 of  $H_f^2$

Maximal values of  $H_f = 2.5 \cdot 10^{16} \text{G}$

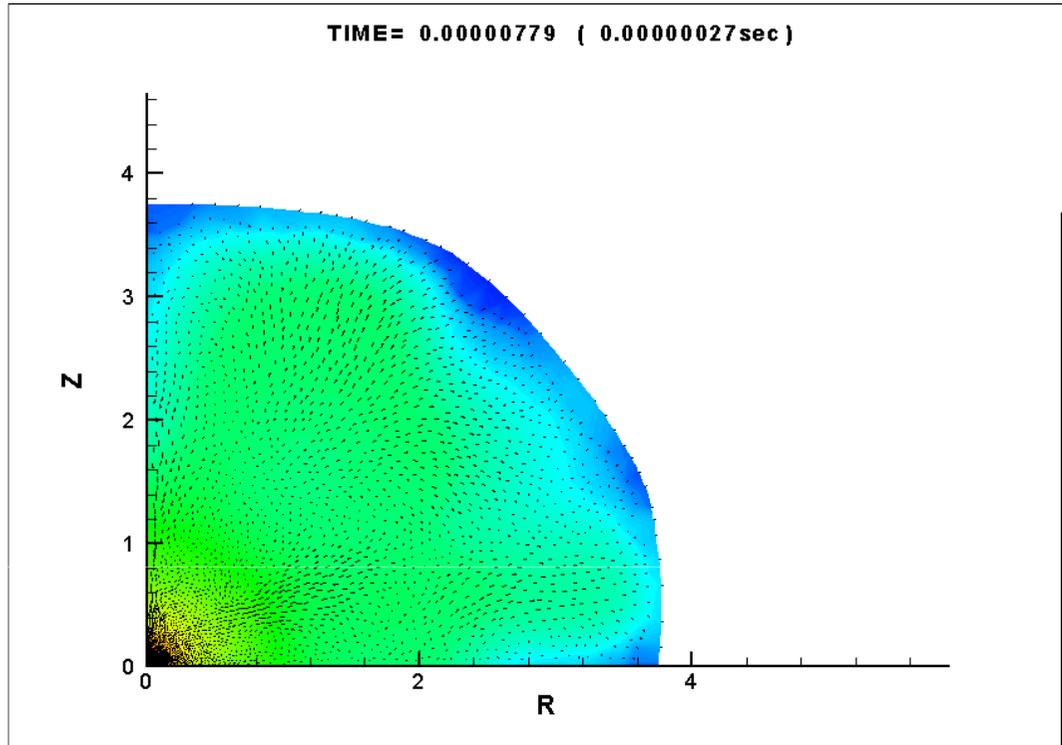
TIME= 0.00000779 ( 0.00000027sec)



The magnetic field  
at the surface of the  
neutron star after the  
explosion is

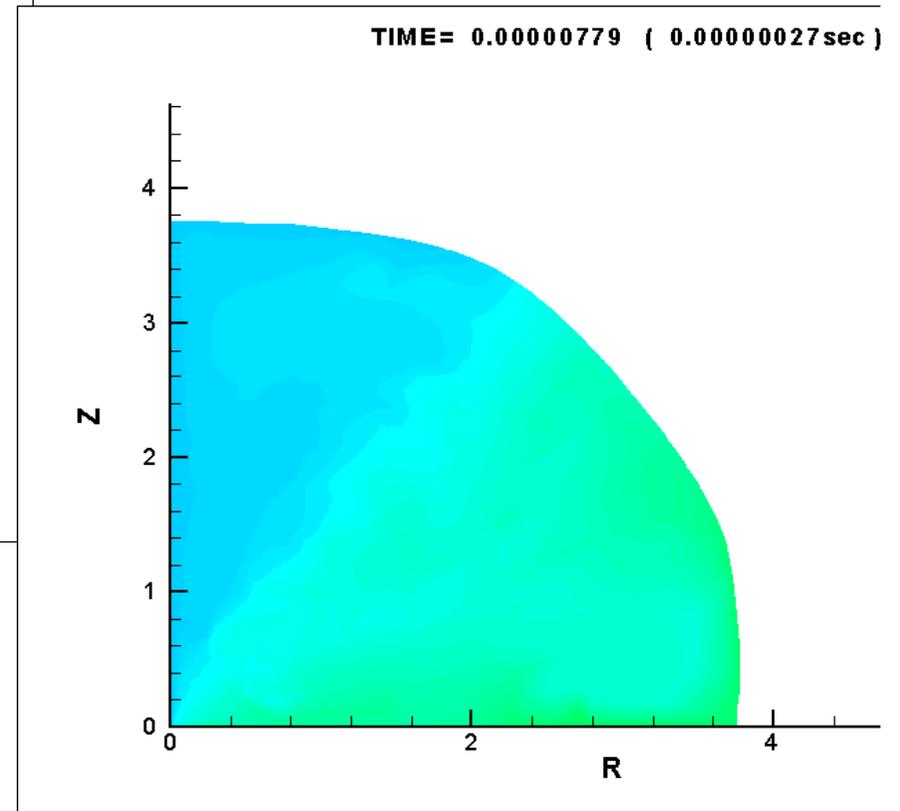
$$H = 4 \cdot 10^{12} \text{Gs}$$

# Temperature and velocity field



# Specific angular momentum

$$rV_{\phi}$$

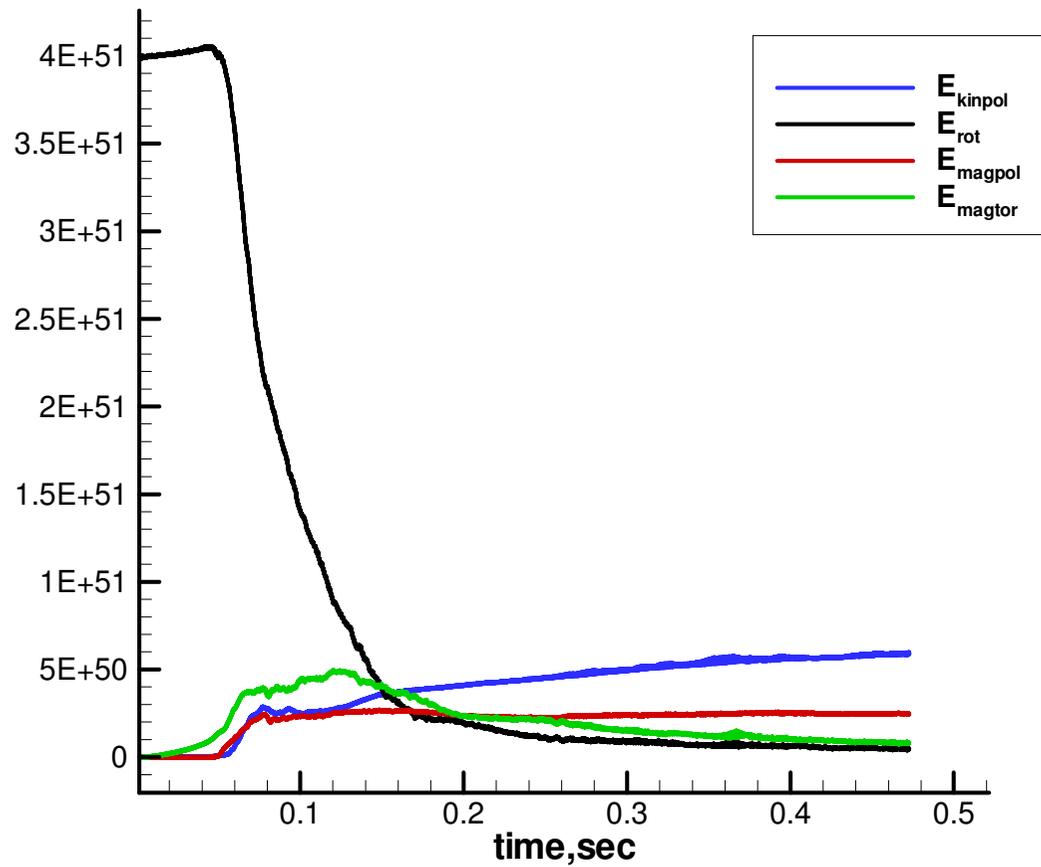


## Rotational energy

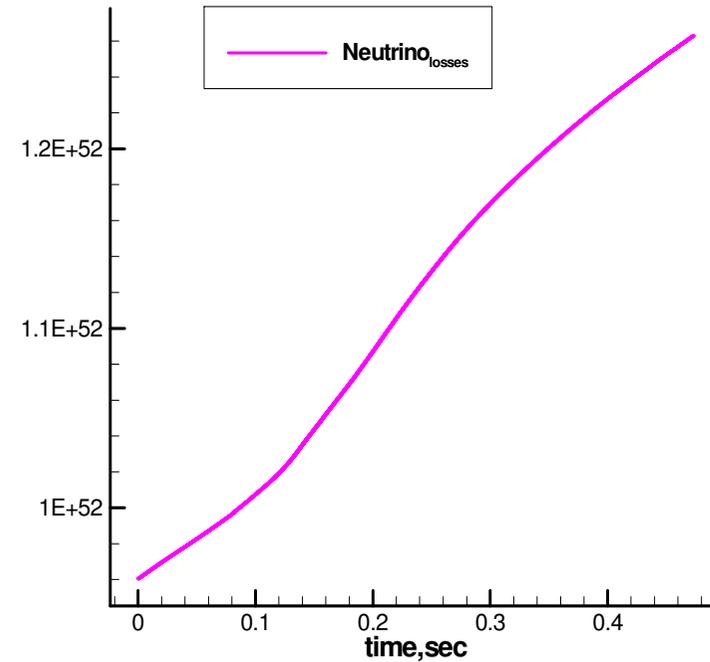
Magnetic poloidal energy

Magnetic toroidal energy

Kinetic poloidal energy



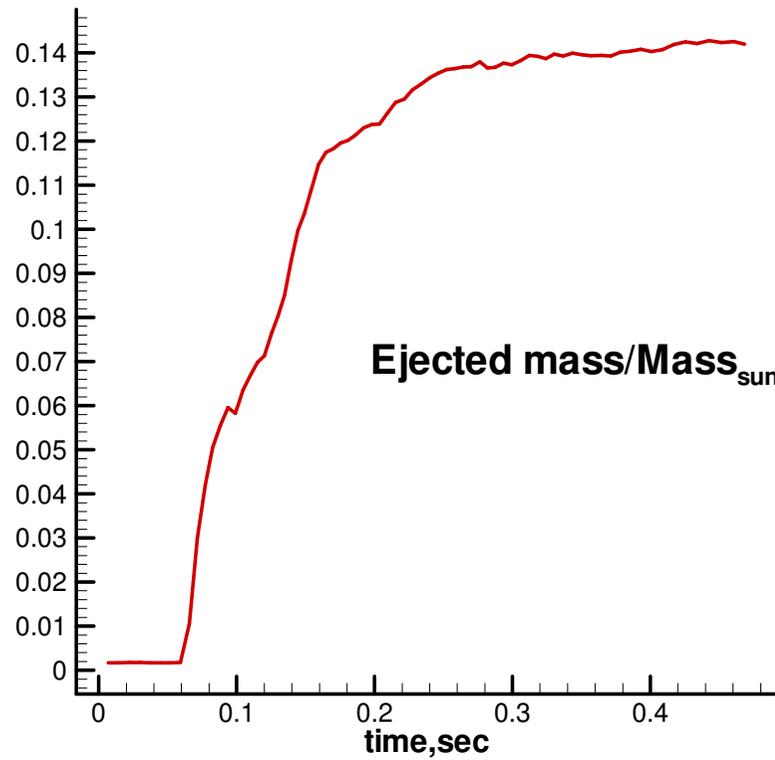
## Neutrino losses



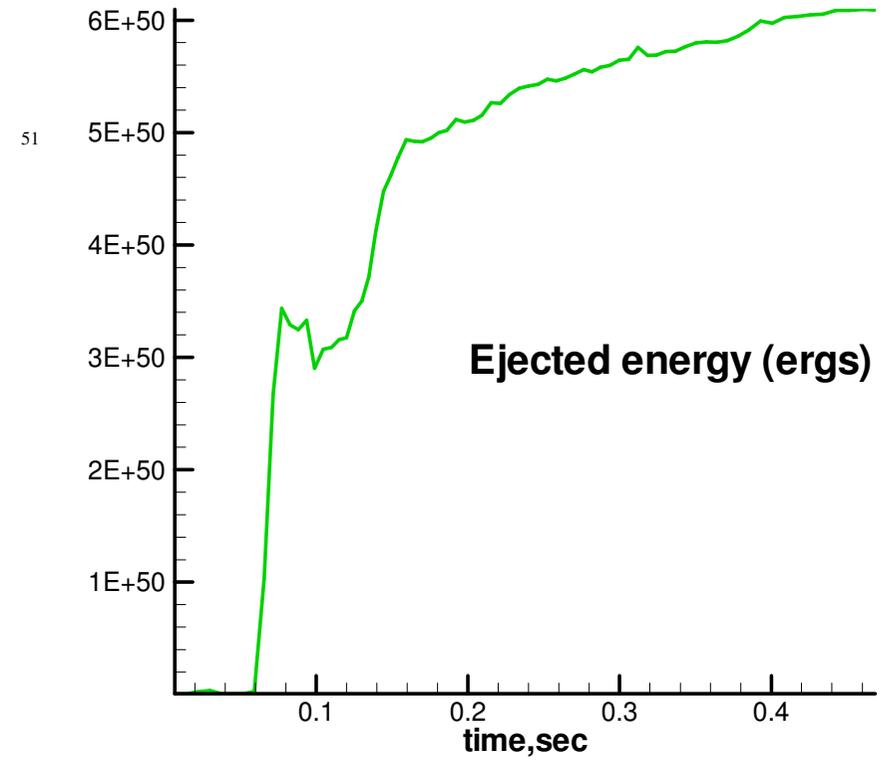
Particle is considered “ejected” if its kinetic energy is greater than its potential energy ( $\alpha=10^{-6}$ )

Ejected energy

$0.6 \cdot 10^{51}$  эрг



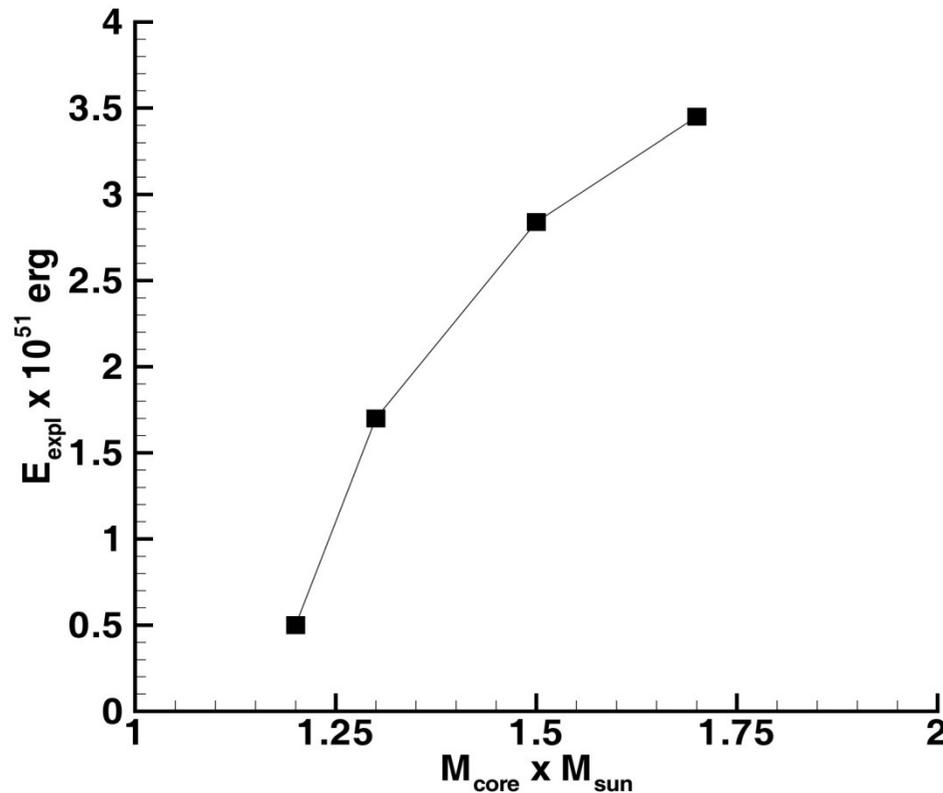
Ejected mass  $0.14M_{\odot}$



# MR supernova – different core masses

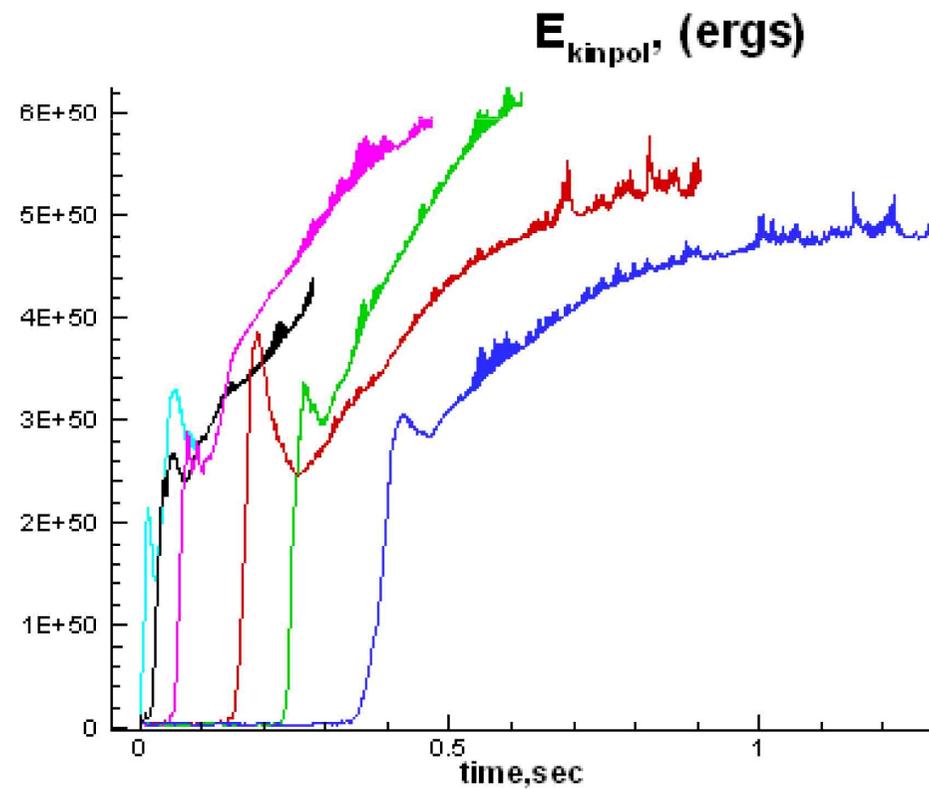
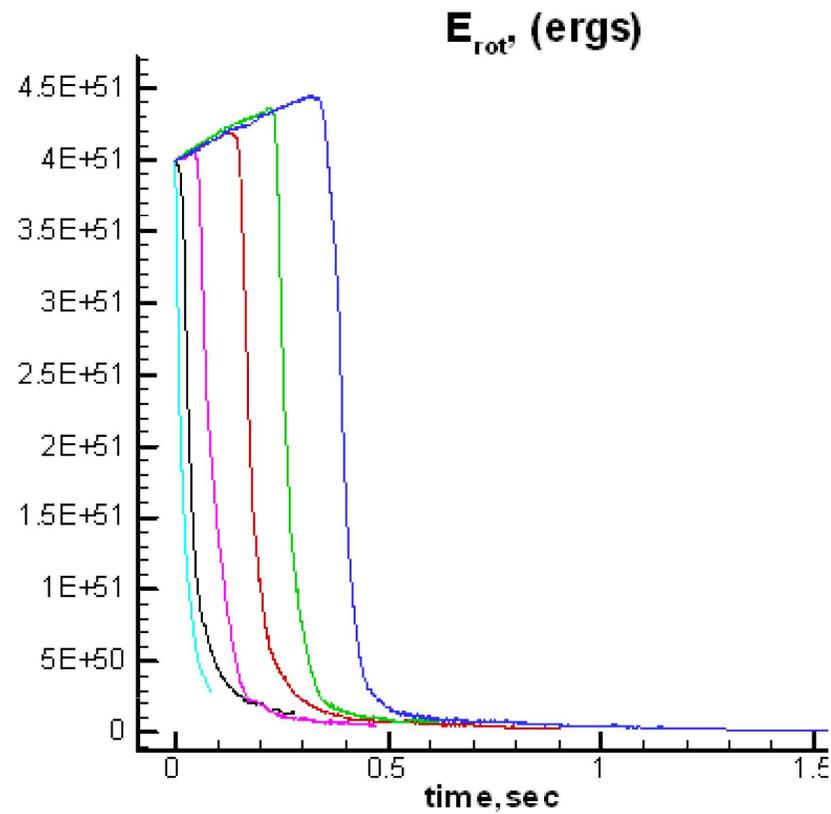
BK, SM, NA, *Astron. Zh.* (2008), 85, 1109

Dependence of the MR supernova explosion energy on the core mass



## Magnetorotational explosion at different

$$\alpha = 10^{-2} - 10^{-12}$$



Dependence of the explosion time on

$$\alpha = \frac{E_{\text{mag}0}}{E_{\text{grav}0}}$$

*1-D calculations:  
Explosion time*

$$t_{\text{взрыва}} \sim \frac{1}{\sqrt{\alpha}}$$

$$\alpha = 10^{-2} \Rightarrow t_{\text{explosion}} = 10,$$
$$\alpha = 10^{-12} \Rightarrow t_{\text{explosion}} = 10^6 (!)$$

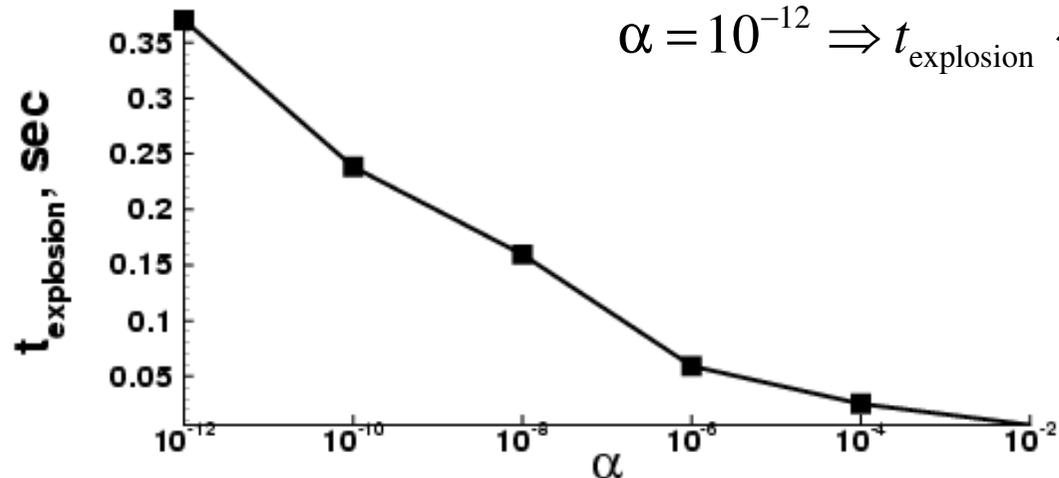
*2-D calculations:  
Explosion time*

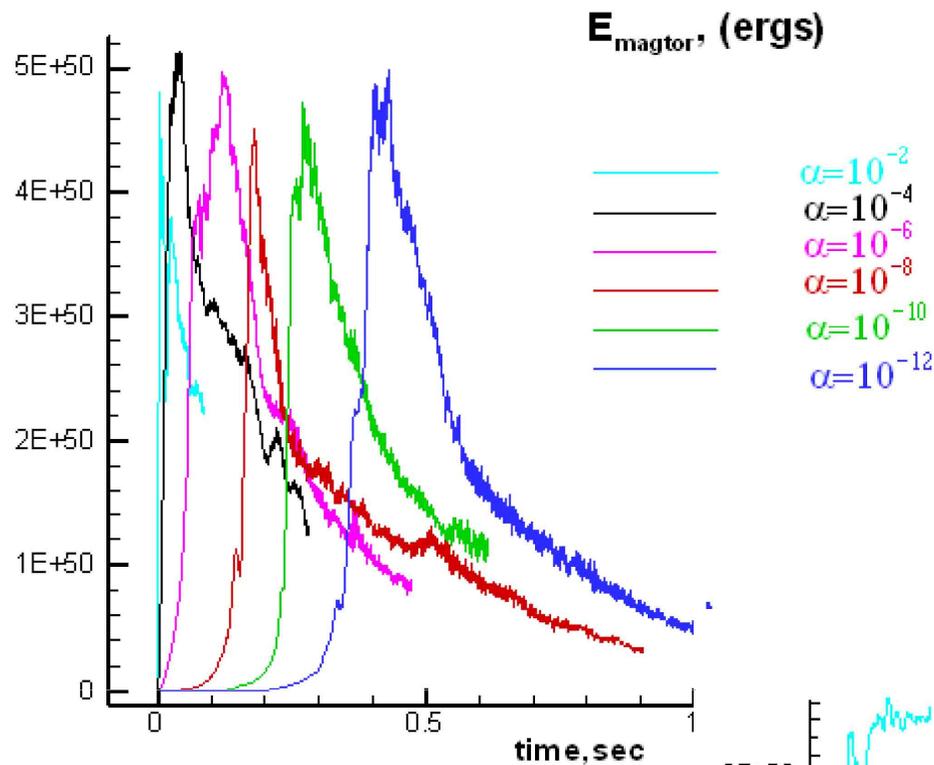
$$t_{\text{взрыва}} \sim -\log(\alpha)$$

(for small  $\alpha$ )

$$\alpha = 10^{-6} \Rightarrow t_{\text{explosion}} \sim 6,$$

$$\alpha = 10^{-12} \Rightarrow t_{\text{explosion}} \sim 12.$$

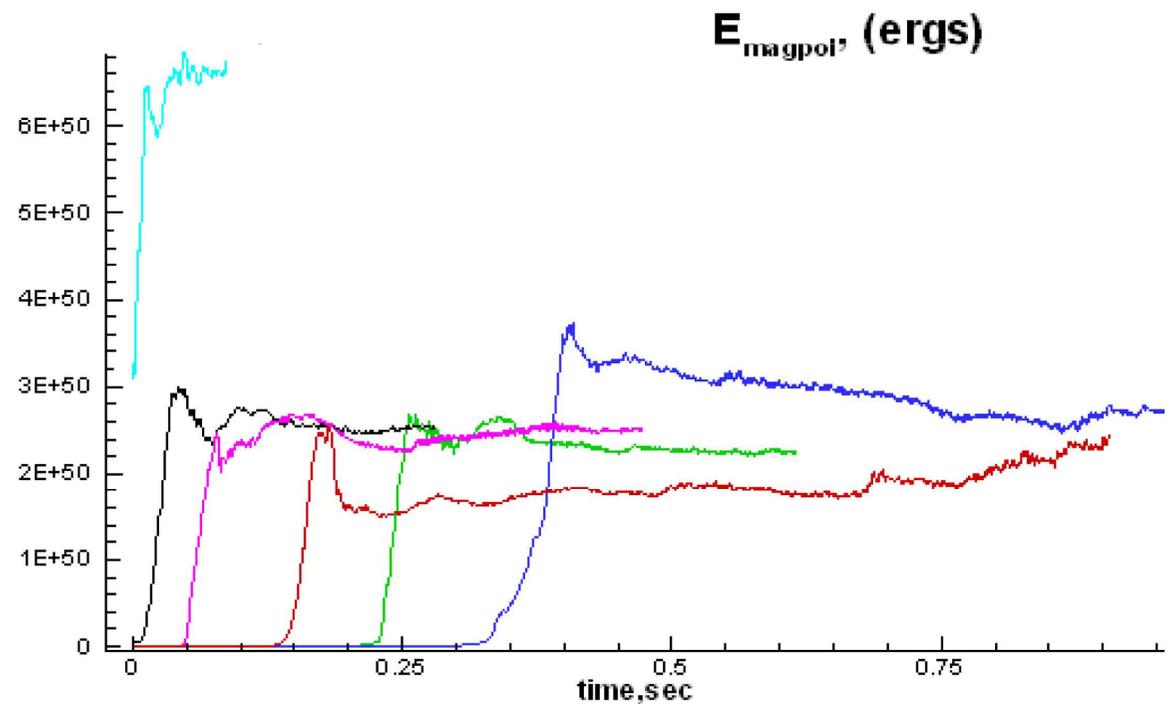




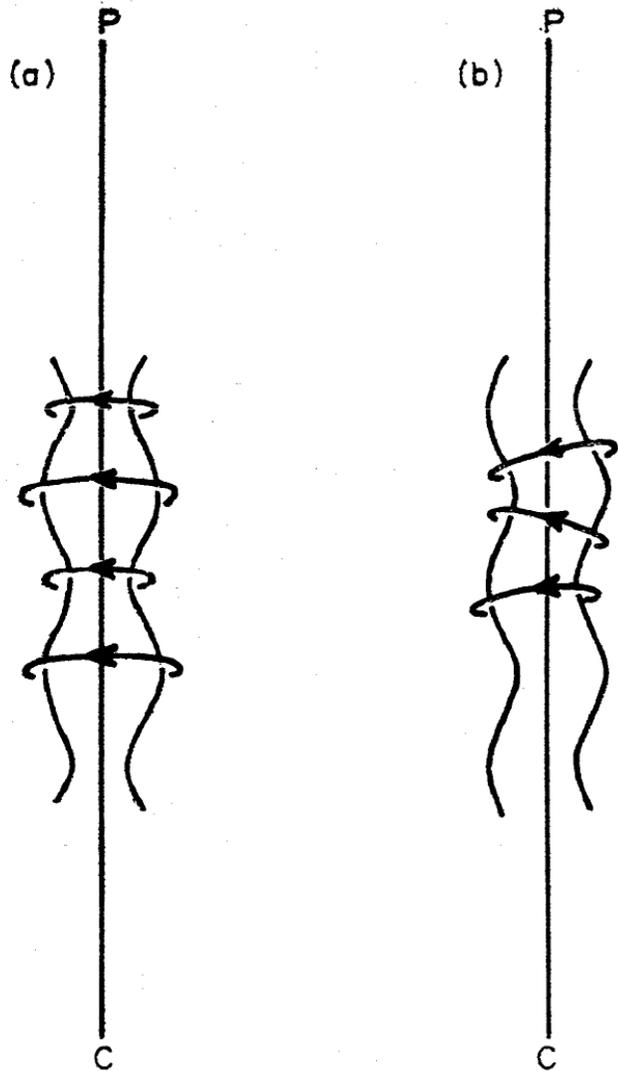
# Magnetorotational instability

$\Rightarrow$  exponential growth of magnetic fields.

Dungey 1958, Velikhov 1959, Spruit 2002, Akiyama et al. 2003



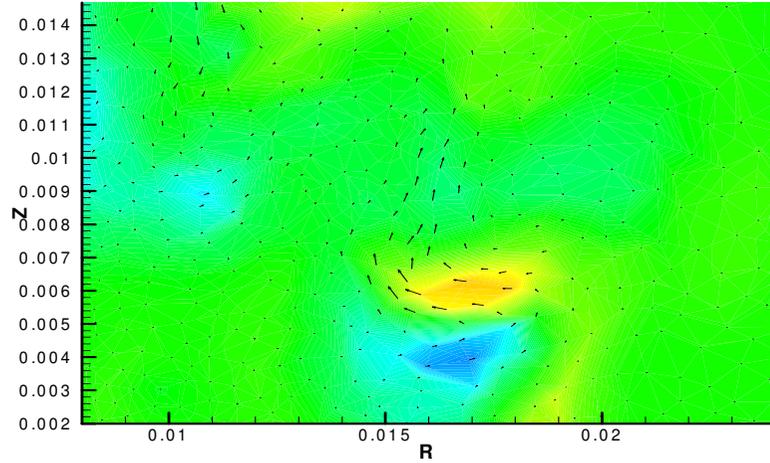
MHD instability: in 2D and 3D – Tayler instability (no rotation)  
Magnetic Differential Rotation Instability (MDRI)



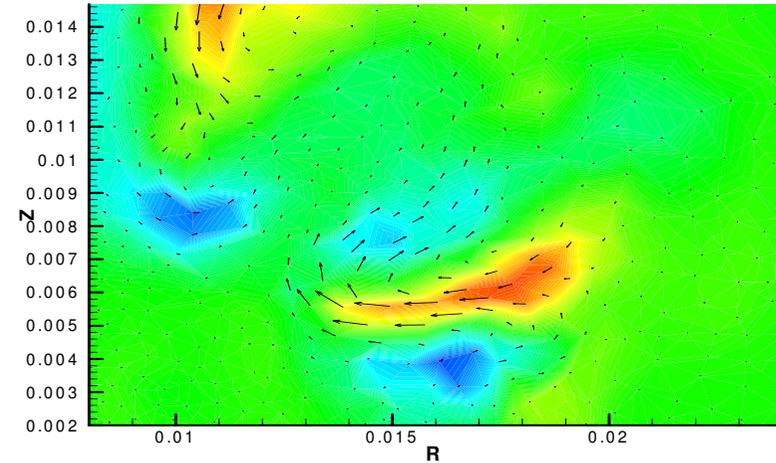
Tayler, R. MNRAS 1973

# Inner region: development of MDRI

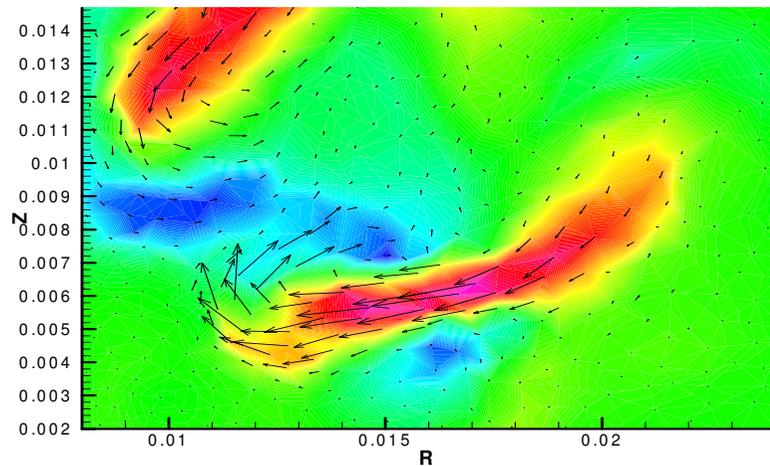
TIME= 34.83616590 ( 1.20326837sec )



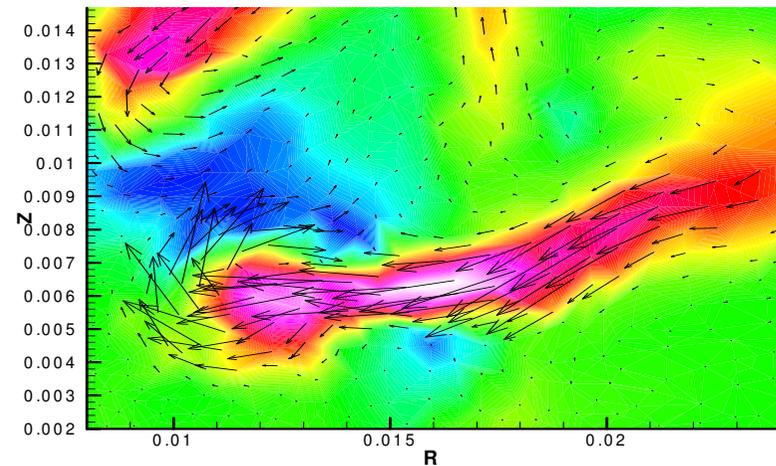
TIME= 35.08302173 ( 1.21179496sec )



TIME= 35.26651529 ( 1.21813298sec )



TIME= 35.38772425 ( 1.22231963sec )



Toroidal (color) and poloidal (arrows) magnetic fields  
(quadrupole)

## Toy model of the MDRI development: exponential growth of the magnetic fields

$$\frac{dH_\varphi}{dt} = H_r \left( r \frac{d\Omega}{dr} \right); \quad \text{at initial stages} \quad H_\varphi < H_\varphi^* : \left( r \frac{d\Omega}{dr} \right) = A \approx \text{const},$$

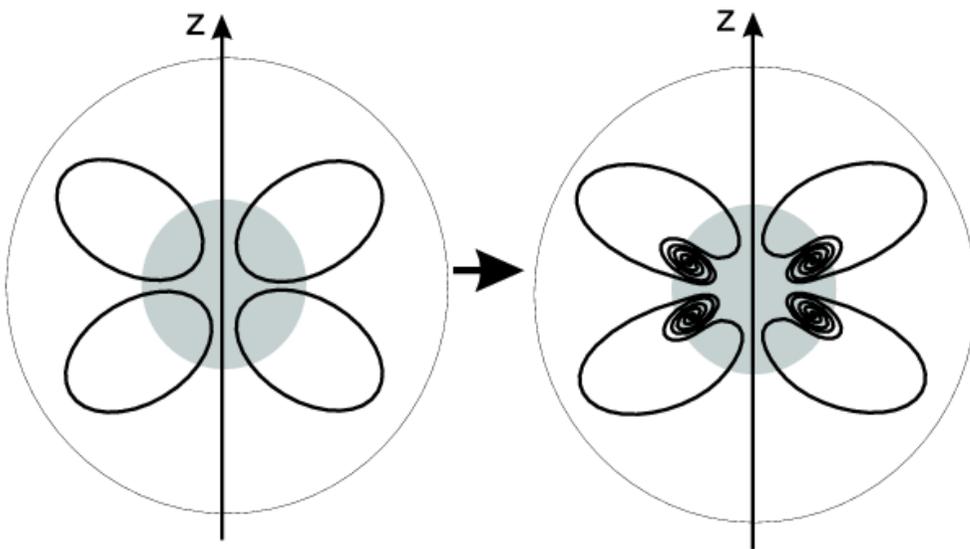
MRI leads to formation of multiple *poloidal* differentially rotating vortices. Angular velocity of vortices is growing (linearly) with a growth of  $H_\varphi$ .

$$\frac{dH_r}{dt} = H_{r0} \left( \frac{d\omega_v}{dl} l \right), \quad \left( \frac{d\omega_v}{dl} l \right) \approx \alpha (H_\varphi - H_\varphi^*).$$

$$\frac{d^2}{dt^2} (H_\varphi - H_\varphi^*) = AH_{r0} \alpha (H_\varphi - H_\varphi^*)$$

⇓

$$\begin{cases} H_\varphi = H_\varphi^* + H_{r0} e^{\sqrt{A\alpha H_{r0}}(t-t^*)}, \\ H_r = H_{r0} + \frac{H_{r0}^{3/2} \alpha^{1/2}}{\sqrt{A}} \left( e^{\sqrt{A\alpha H_{r0}}(t-t^*)} - 1 \right). \end{cases}$$

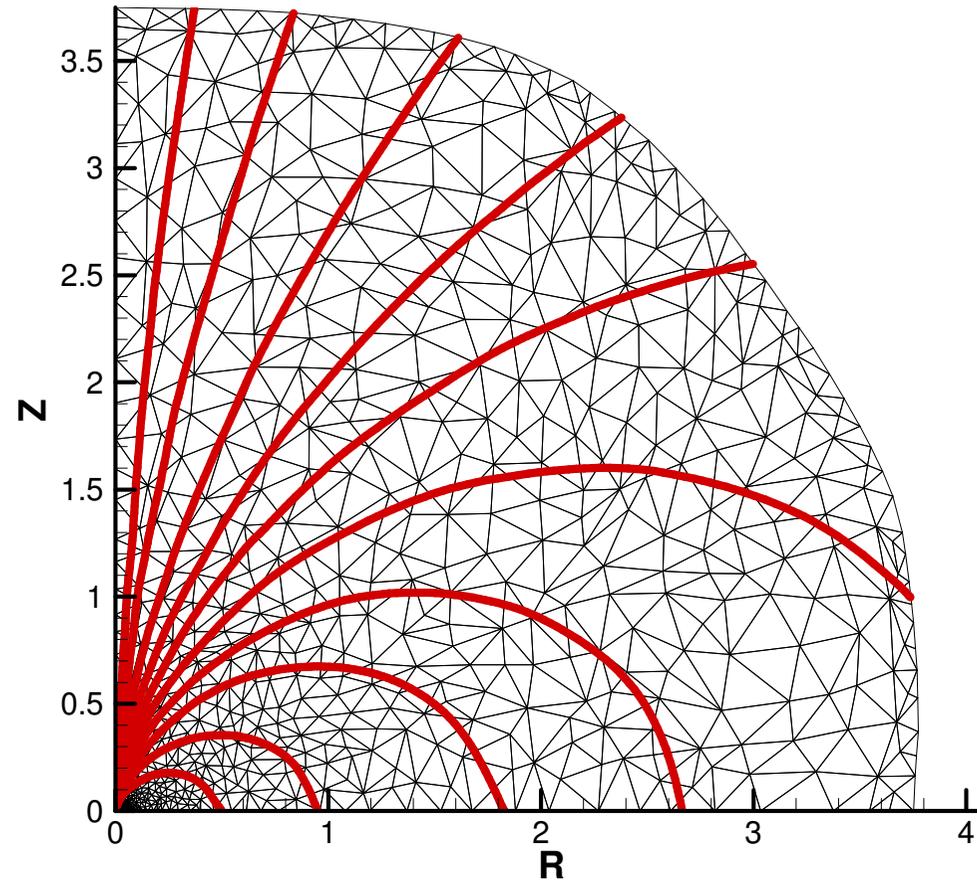


# Jet formation in MDRE

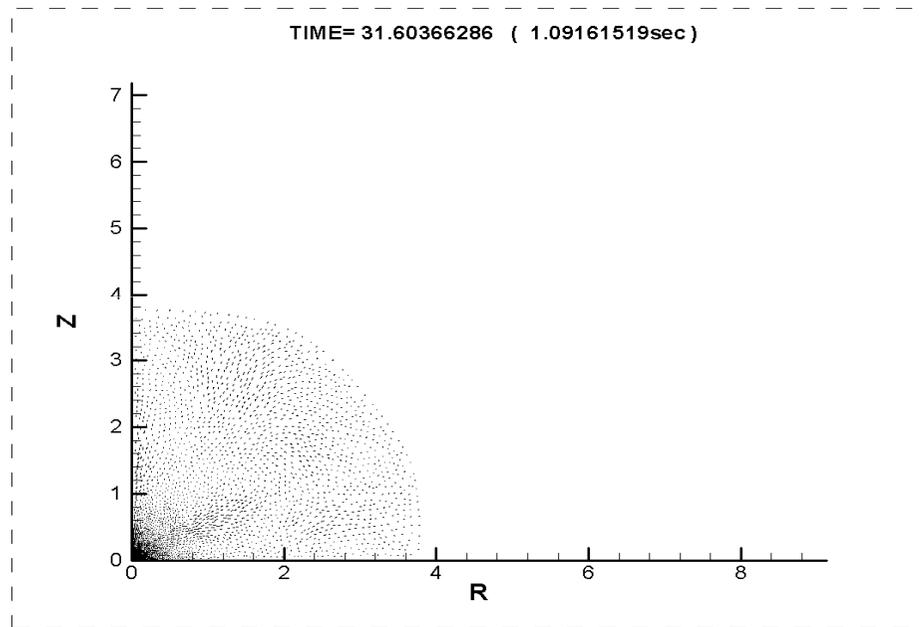
**Moiseenko et al.** Astro-ph/0603789

Magnetorotational supernovae with jets

**Dipole-like** initial magnetic field



## Jet formation in MDRE: velocity field evolution

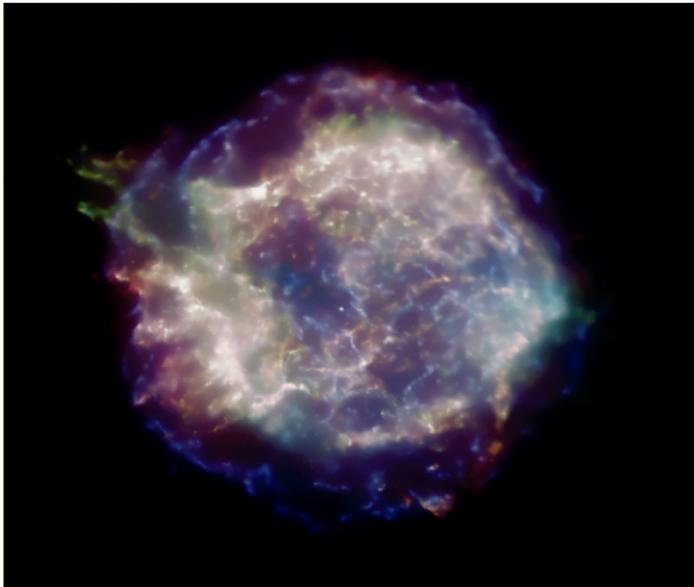




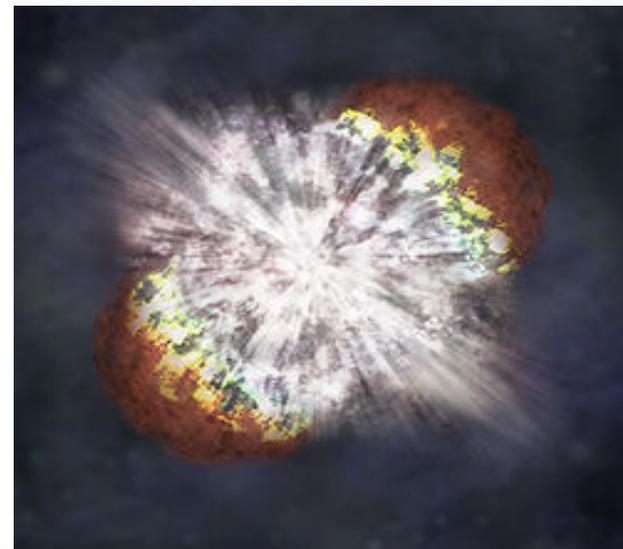
Composite Optical/X-ray image of the **Crab Nebula**, showing **ynchrotron emission** in the surrounding **pulsar wind nebula**, powered by injection of magnetic fields and particles from the central pulsar.



The Vela Pulsar and its surrounding pulsar wind nebula.

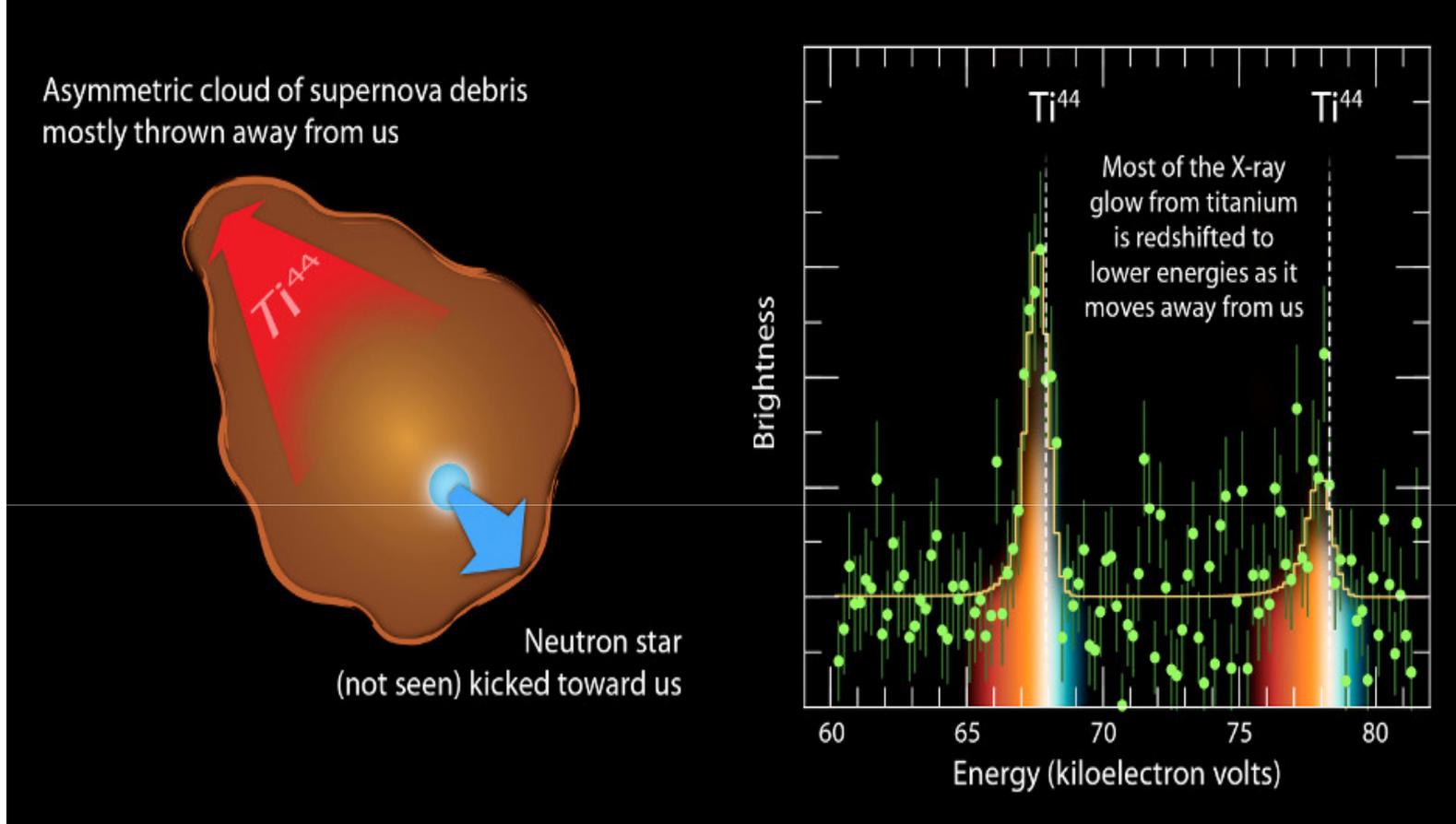


**X-ray images of Cassiopeia A, taken by the Chandra satellite**



NASA's artist impression of SN 2006gy, one of the most luminous hypernovae seen

## NuSTAR Sees Titanium Glow in Supernova 1987A

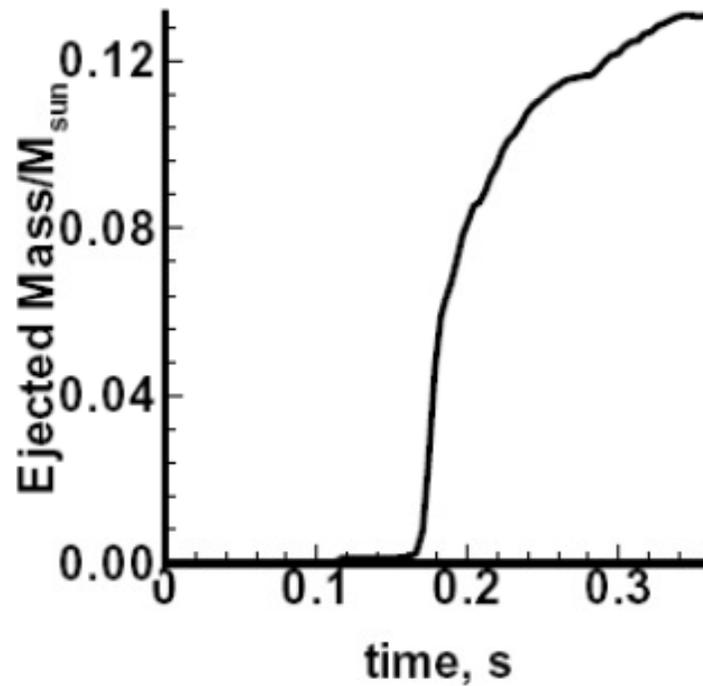
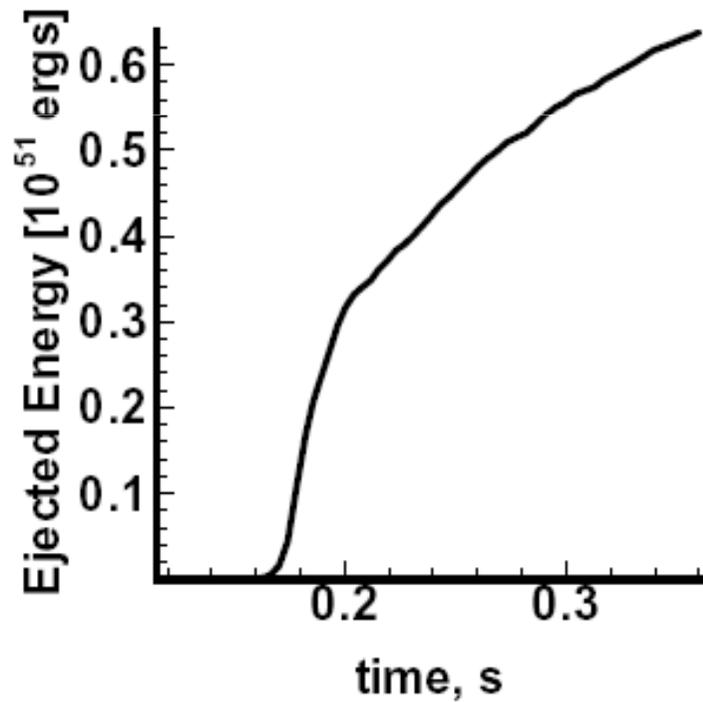


The plot of data from NASA's Nuclear Spectroscopic Telescope Array, or NuSTAR (right), amounts to a "smoking gun" of evidence in the mystery of how massive stars explode. The observations indicate that supernovae belonging to a class called Type II or core-collapse blast apart in a lopsided fashion, with the core of the star hurtling in one direction, and the ejected material mostly expanding the other way

Jet formation in MDRE: (**dipole magnetic field**)

Energy of explosion  $\approx 0.6 \cdot 10^{51} \text{ erg}$

Ejected mass  $\approx 0.14 M_{\odot}$



## RECENT RESULTS

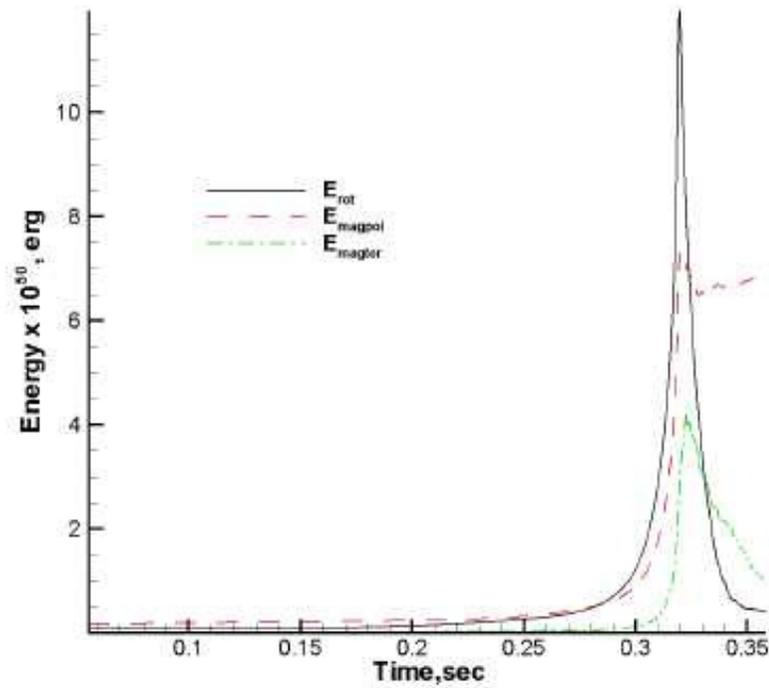
Improvement of EoS, neutrino emission rate, neutrino transfer description (**Shen et al., 1998; K. Sato et al. 2005-2010**)

Calculations show results similar to what was obtained using a simplified description of physical processes.

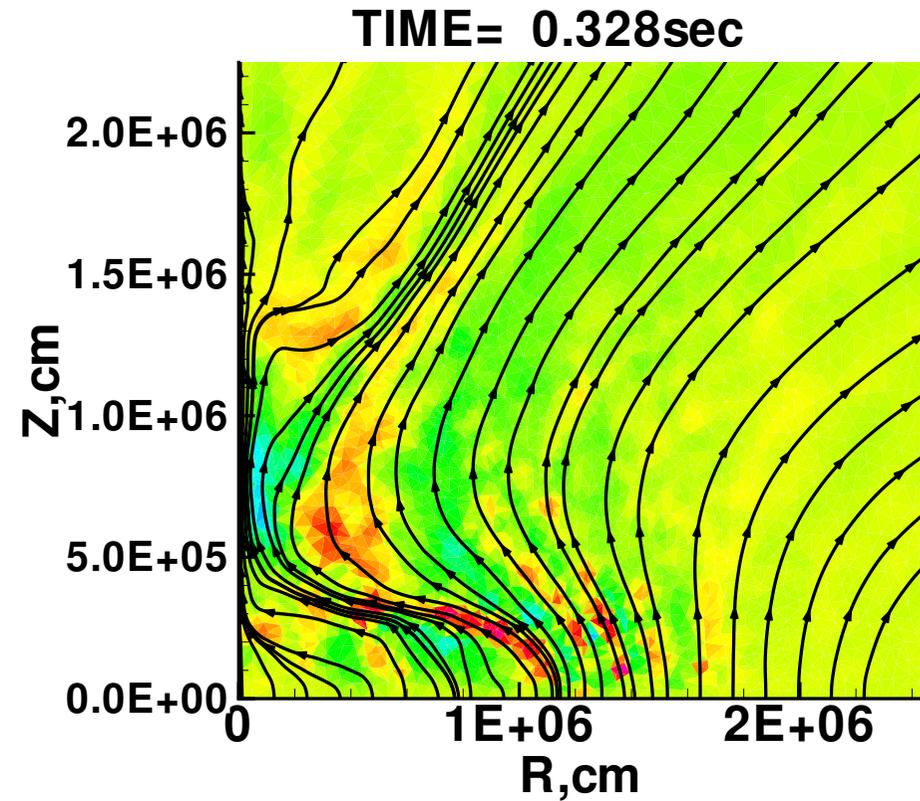
**Bisnovatyi-Kogan, G. S.; Moiseenko, S. G.; Ardeljan, N. V.**  
Numerical Modeling of Space Plasma Flows (ASTRONUM2012).  
Proceedings of a 7th International Conference. ASPC, 474, 47 (2013)

**S.G.Moiseenko, G.S.Bisnovatyi-Kogan, K.Kotake, T.Takiwaki,**  
**K.Sato** (in preparation)

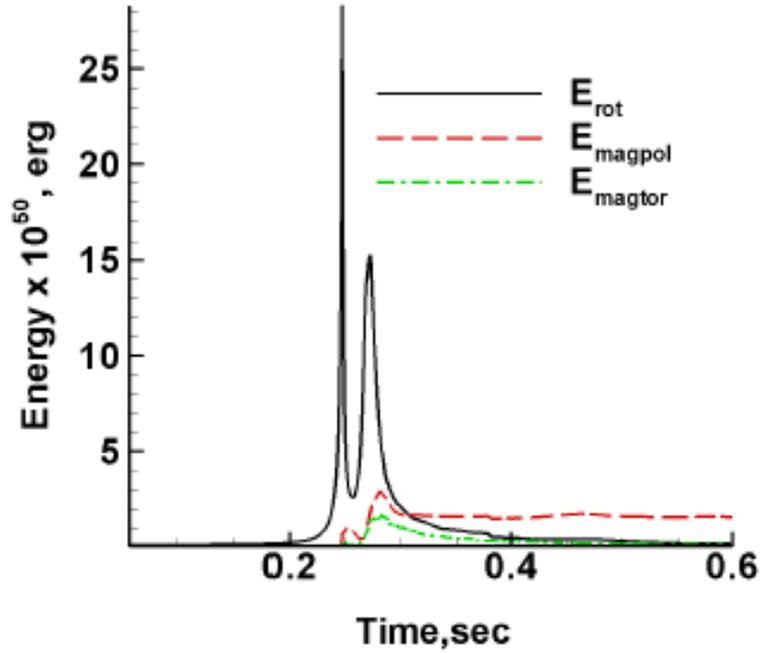
**S.G.Moiseenko, G.S.Bisnovatyi-Kogan**  
**Development of the Magneto-Differential-Rotational Instability**  
**in Magnetorotational Supernova**  
*Astronomy Reports, 2015, Vol. 59, No. 7, pp. 573–580.*



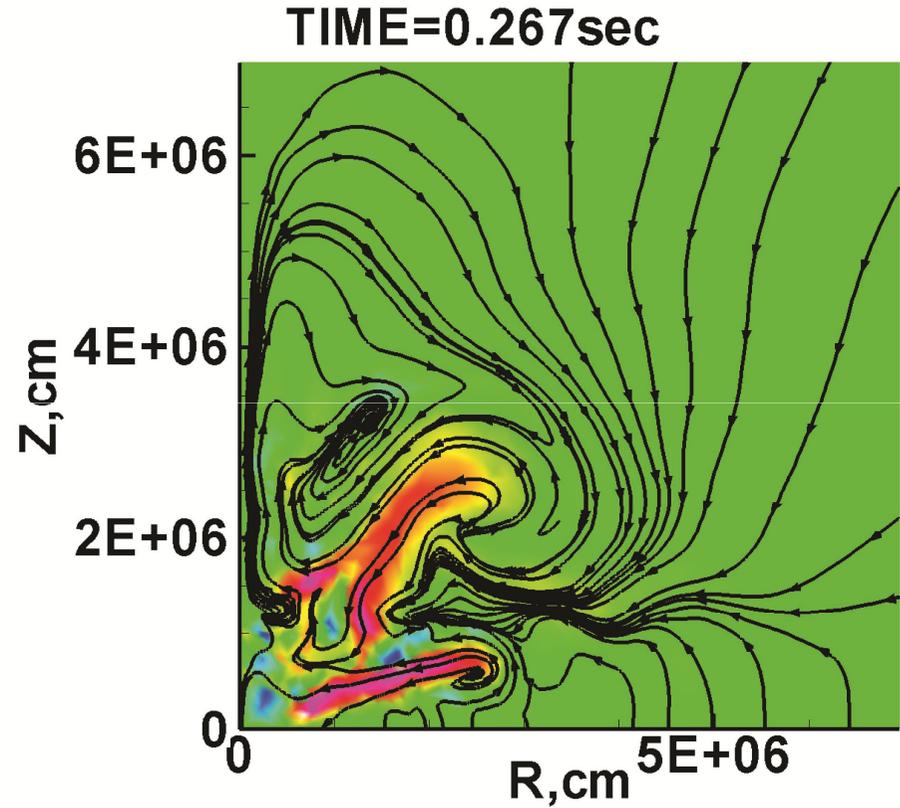
**Figure 7.** Time evolution of rotational energy  $E_{rot}$  (solid line), magnetic poloidal energy  $E_{magpol}$  (dashed line) and magnetic toroidal energy  $E_{magtor}$  (dash-dotted line) for the case  $H_0 = 10^{12}$ , G.,  $E_{rot0}/E_{grav0} = 1\%$ .



Absence of MDRI at  $t = 328$  ms for the case  $H_0 = 10^{12}$  G,  $E_{rot0}/E_{grav0} = 1\%$  (contour plot - the toroidal magnetic field, black lines indicate the poloidal magnetic field).

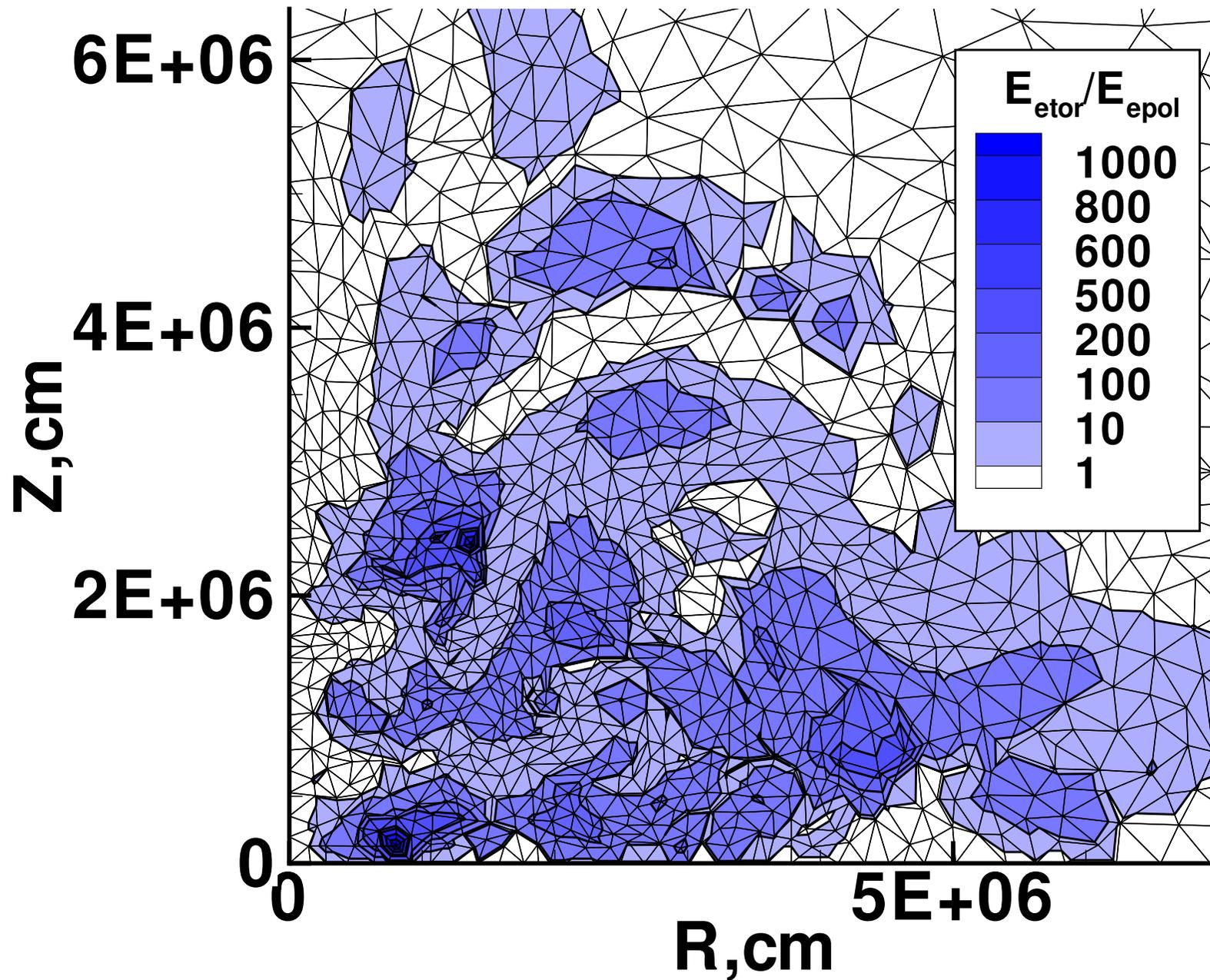


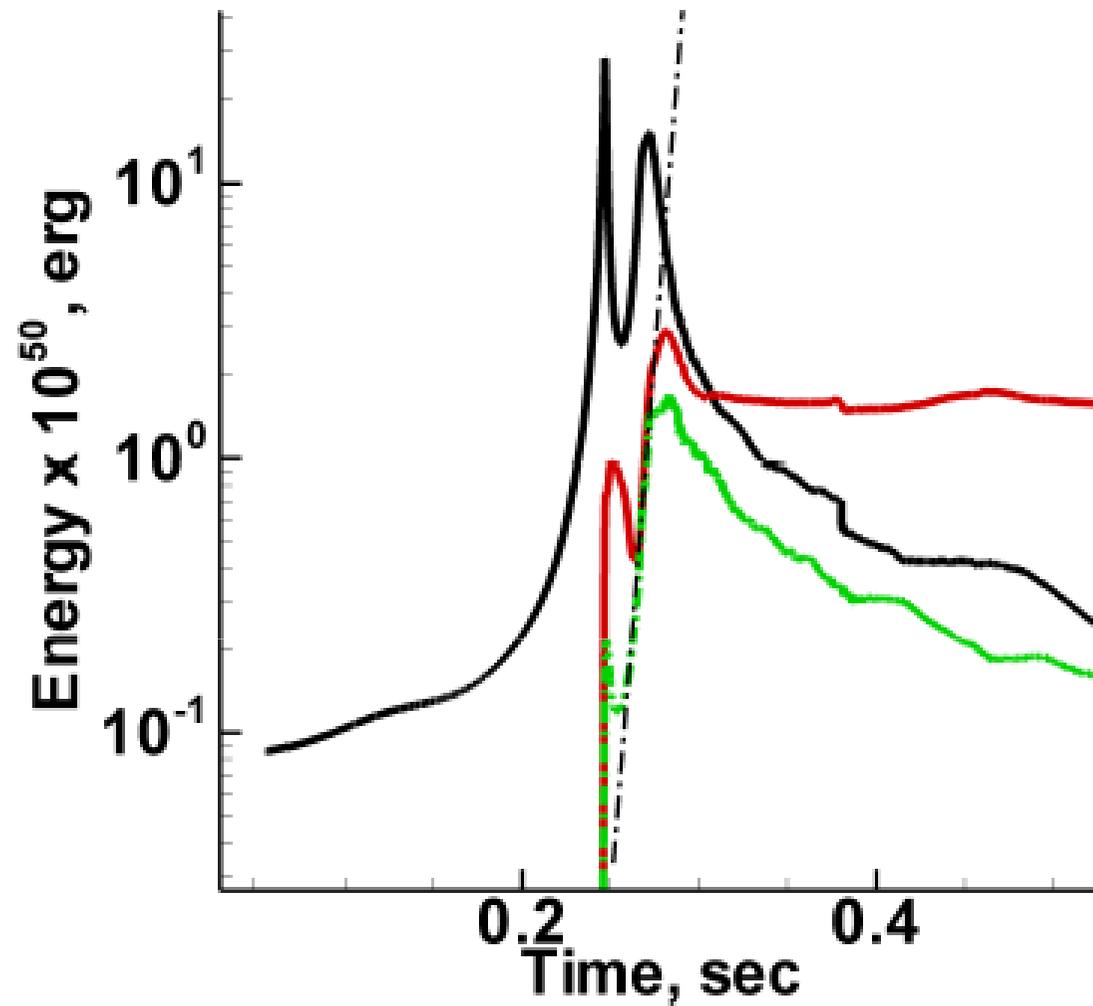
**Figure 4.** Time evolution of rotational energy  $E_{rot}$  (solid line), magnetic poloidal energy  $E_{magpol}$  (dashed line) and magnetic toroidal energy  $E_{magtor}$  (dash-dotted line) for the case  $H_0 = 10^9$  G,  $E_{rot0}/E_{grav0} = 1\%$ .



**Figure 2.** Developed MDRI due to convection and MRI/Taylor instability at  $t = 267$  ms for the case  $H_0 = 10^9$  G,  $E_{rot0}/E_{grav0} = 1\%$  (contour plot - the toroidal magnetic field, arrow lines - force lines of the poloidal magnetic field).

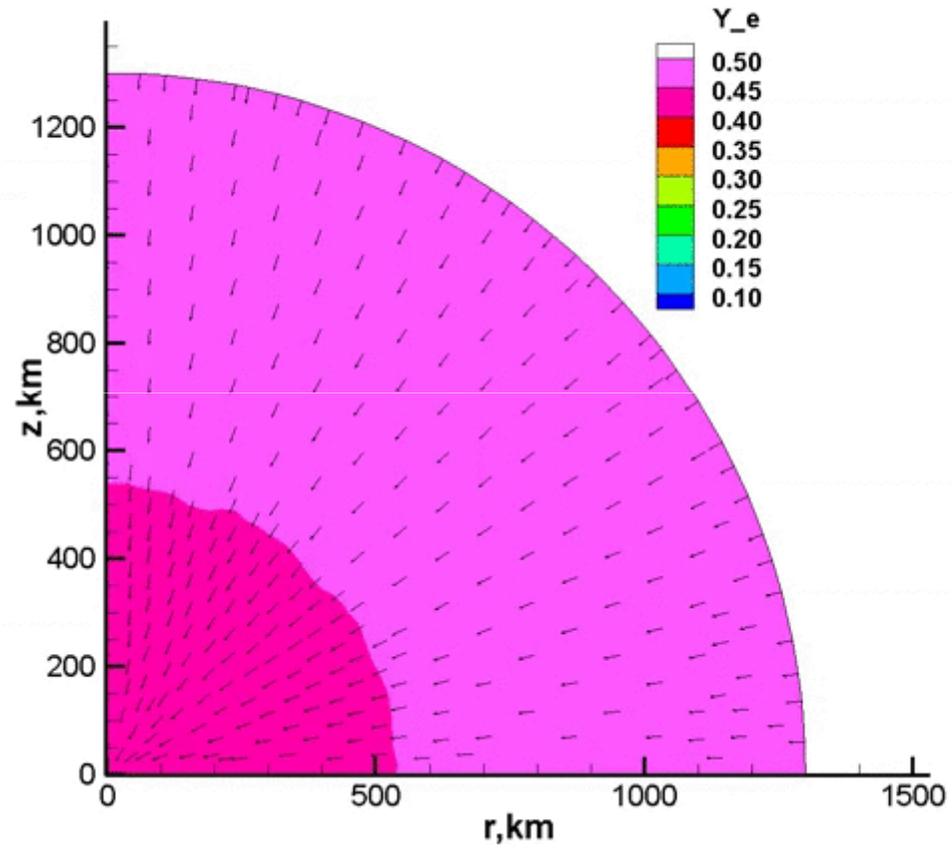
**TIME= 0.265sec**





**Figure 5.** Zoomed time evolution of rotational energy  $E_{rot}$  (solid line), magnetic poloidal energy  $E_{magpol}$  (dashed line) and magnetic toroidal energy  $E_{magtor}$  (dash-dotted line) for the case  $H_0 = 10^9 \text{G}$ ,  $E_{rot0}/E_{grav0} = 1\%$ . Straight dash-dotted line shows exponential growth of the toroidal and poloidal magnetic energies.

TIME= 0.00001000 ( 0.00000035sec )



# Asymmetry of the explosion

CP violation in weak processes in regular magnetic field: does not work,  
because MRI leads to formation of highly chaotic field.

Astro-ph/0510229

MULTI-DIMENSIONAL RADIATION HYDRODYNAMIC SIMULATIONS OF  
PROTONEUTRON STAR CONVECTION

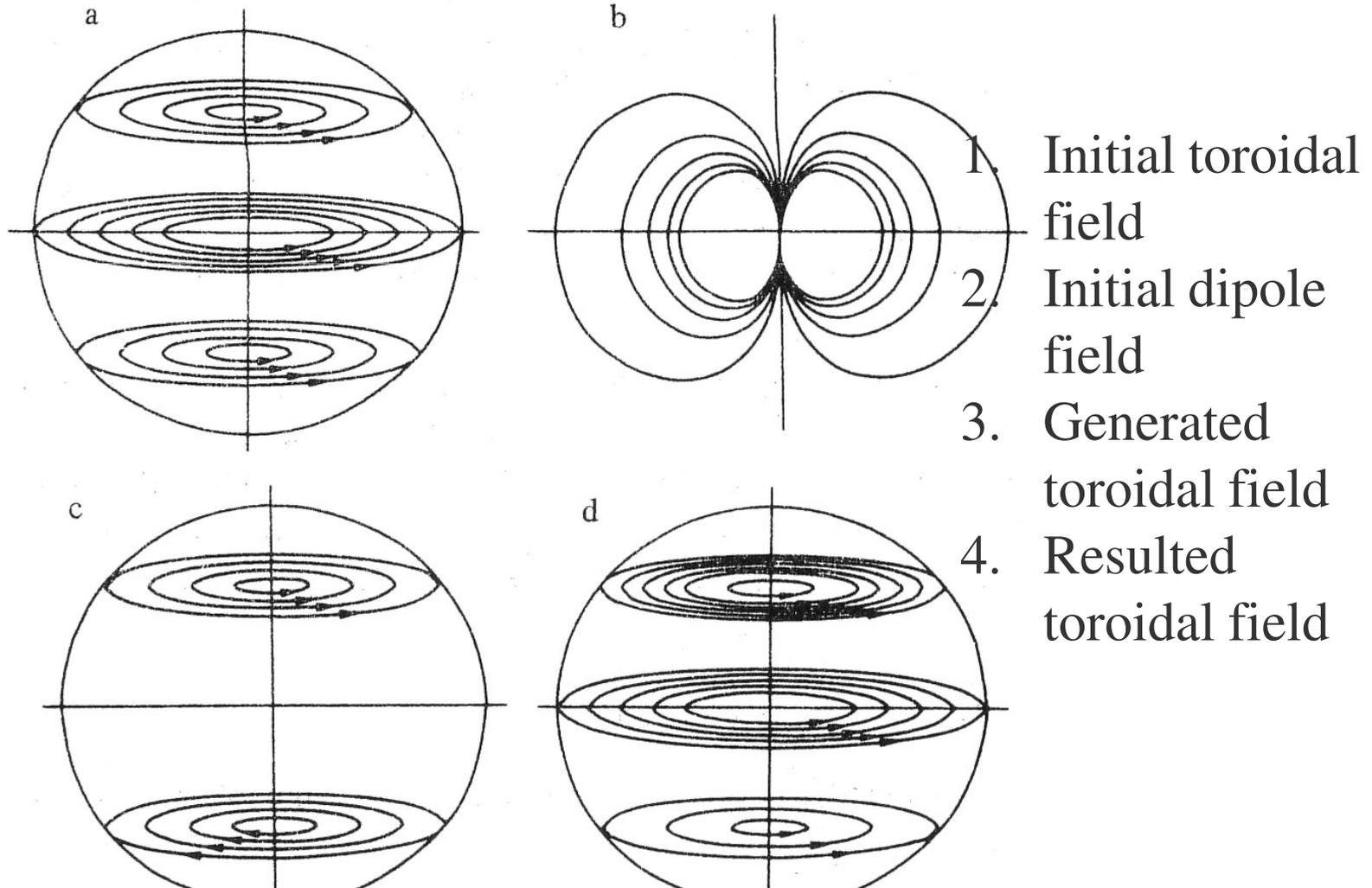
L. Dessart, A. Burrows, E. Livne, C.D. Ott

PNS convection is thus found to be a secondary feature  
of the core-collapse phenomenon, rather than a  
decisive ingredient for a successful explosion.

# Asymmetry of the explosion

## Violation of mirror symmetry of magnetic field

(BK, Moiseenko, 1992 *Astron. Zh.*, 69, 563 (SvA,  
1992, 36, 285)



BK, 1993, Astron. Ap. Transactions **3**, 287

## Interaction of the neutrino with asymmetric magnetic field

$$B_c = \frac{m_e^2 c^3}{e \hbar} = 4.4 \times 10^{13} \text{Gs}$$

The probability of the neutron decay  $W_n$  in vacuum is

$$W_n = W_0 [1 + 0.17(B/B_c)^2 + \dots] \quad \text{at } B \ll B_c$$

$$W_n = 0.77W_0(B/B_c) \quad \text{at } B \gg B_c$$

Dependence of the weak interaction cross-section on the magnetic field strength lead to the asymmetric neutrino flux and formation of rapidly moving pulsars due to the recoil action as well as rapidly moving black holes.

Neutrino heat conductivity

$$H_\nu = -\frac{7}{8} \frac{4acT^3}{3} l_T \frac{\partial T}{\partial r} \quad \kappa_\nu = 1/(l_T \rho)$$

energy flux neutrino opacity

**The anisotropy of the flux**

$$\delta_L = \frac{L_+ - L_-}{L_+ + L_-}$$

**Kick velocity along the rotational axis**

## Approximate estimation

$$v_{nf} = \frac{2}{\pi} \frac{L_\nu}{M_n c} \frac{P B_{\phi 0}}{|B_p|} \left( 0.5 + \ln \left( \frac{20 \text{ s}}{P} \frac{|B_p|}{B_{\phi 0}} \right) \right)$$

$x = \frac{B_{\phi 0}}{|B_p|}$  between 20 and  $10^3$ , we have  $v_{nf}$  between 140 and 3000 km/s

Important to do:

Numerical simulations without mirror symmetry

Accurate formulae for neutrino processes at high magnetic field

T. V. SMIRNOVA, V. I. SHISHOV, AND V. M. MALOFEEV

THE SPATIAL STRUCTURE OF PULSAR EMISSION SOURCES DETERMINED USING  
INTERSTELLAR SCINTILLATION

THE ASTROPHYSICAL JOURNAL, 462:289–295, 1996 May 1

we can conclude that the velocity vector of the movement of the pulsar in space is close to the orientation of the rotation axis of the pulsar. Of course, this interpretation assumes

S. Johnston et al. astro/ph 0510260 (MNRAS, 2005, **364**, 1397)

### **Evidence for alignment of the rotation and velocity vectors in pulsars**

We present strong observational evidence for a relationship between the direction of a pulsar's motion and its rotation axis. We show carefully calibrated polarization data for 25 pulsars, 20 of which display linearly polarized emission from the pulse longitude at closest approach to the magnetic pole... **we conclude that the velocity vector and the rotation axis are aligned at birth.**

S. Johnston, M. Kramer, A. Karastergiou, G. Hobbs, S. Ord , J. Wallman  
MNRAS, 381, Issue 4, pp. 1625 (2007)

**Evidence for alignment of the rotation and velocity vectors in  
pulsars. II. Further data and emission heights**

**Pulsar Spin–Velocity Alignment: Further Results and  
Discussion**

A. Noutsos, M. Kramer, P. Carr and S. Johnston

arXiv:1205.2305v1 [astro-ph.GA] 10 May 2012

We estimate that the observed alignment is robust to within 10%  
systematic uncertainties on the determination of the spin-axis  
direction from polarisation data.

W.H.T. Vlemmings et al. astro-ph/0509025 (Mm. SAI, 2005, **76**, 531)

## **Pulsar Astrometry at the Microarcsecond Level**

Determination of pulsar parallaxes and proper motions addresses fundamental astrophysical questions. We have recently finished a VLBI astrometry project to determine the proper motions and parallaxes of 27 pulsars, thereby doubling the total number of pulsar parallaxes. Here we summarize our astrometric technique and present the discovery of a pulsar moving in excess of 1000 kms, **PSR B1508+55**.

# Conclusions

1. In the magnetorotational explosion (MRE) the efficiency of transformation of rotational energy into the energy of explosion is 10%. This is enough for producing core – collapse SN from rapidly rotating magnetized neutron star.
2. Development of magneto-differential-rotational instability strongly accelerate MRE, at lower values of the initial magnetic fields.
3. The new born neutron star has inside a large (about  $10^{14}$  Gauss) chaotic magnetic field.
4. Jet formation is possible for dipole-like initial topology of the field: possible relation to cosmic gamma-ray bursts; equatorial ejection happens at prevailing of the quadrupole-like component.
5. Development of magneto-differential-rotational instability (MDRI) starts in regions, where the toroidal component significantly prevails the poloidal one.
6. MRE explosion energy is not sensitive to the precision of the input microphysics, and to the numerical scheme.